

Parton Distributions on a Euclidean Lattice---the LaMET Method

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Parton Model

- Leading twist t ($t = 2$) contribution in QCD
- Factorization

$$\sigma(eP \rightarrow e'X) = \int_0^1 dx f_i(x, \mu) \sigma(ei \rightarrow e'i', \mu)$$

$\Lambda_{QCD} \qquad Q$

- Expansion in powers of Λ_{QCD}/Q

Power correction

$$\left(\frac{\Lambda_{QCD}}{Q} \right)^{t-2}$$

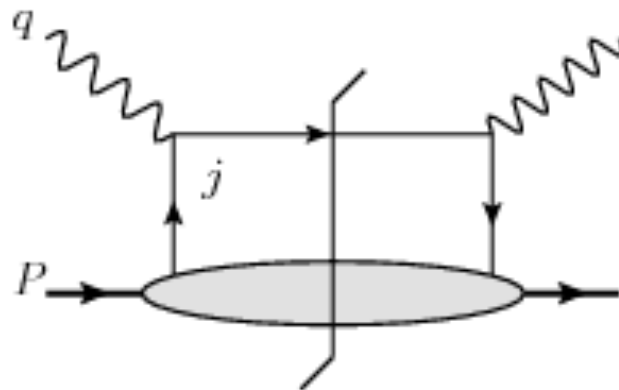
Main result

- Twist-2 matrix element

$$\langle P | \hat{\mathcal{O}}^{\mu_1 \cdots \mu_n}(x) | P \rangle = 2 \langle x^{n-1} \rangle P^{(\mu_1} \cdots P^{\mu_n)}$$

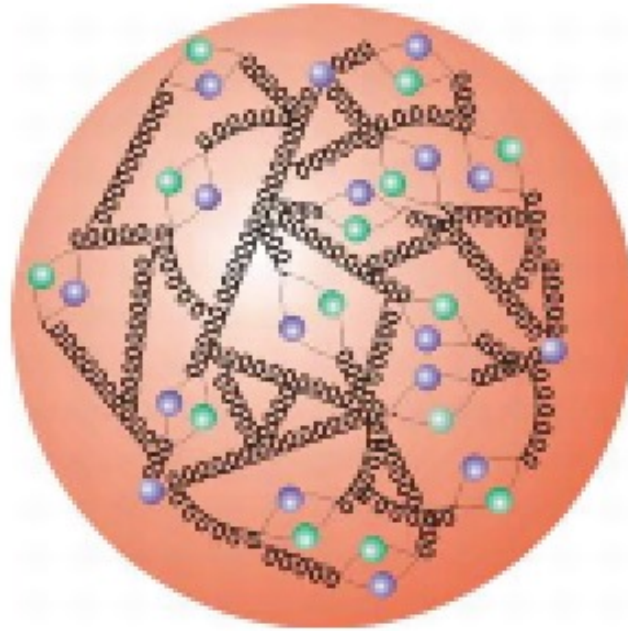
Wilson line

$$q(x, \mu^2) = \int \frac{d\xi^-}{4\pi} e^{ix\xi^- P^+} \langle P | \bar{\psi}(0) \lambda \cdot \gamma \Gamma \psi(\xi^- \lambda) | P \rangle$$



The struck parton moves on a light cone yields a sum of twist-2 operators after OPE.

Golden Age for Proton Structure Studies



Parton structures: 1d mom+spin PDF to 3d GPD & TMD to Wigner (and beyond?) [BNL, JLab, J-PARC, COMPASS, GSI, **EIC**, LHeC, ...] to **applications** (Higgs, new physics...)

PDFs from QCD---Why is it so hard?

- The number of quark anti-quark pairs diverges (manifestation of non-perturbative nature of the problem): .

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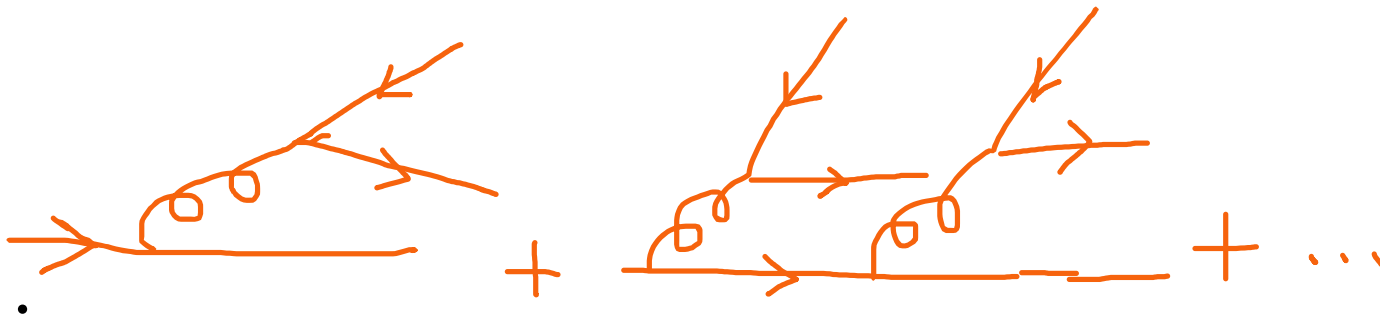
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- Euclidean lattice: light cone operators cannot be distinguished from local operators.

$$\begin{aligned} t^2 - \mathbf{r}^2 &= 0 \\ -t_E^2 - \mathbf{r}^2 &= 0 \end{aligned}$$

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PDF $\propto \delta(x)$. While short distance expansion is natural in Euclidean space, lightcone expansion is not.

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PDFs from QCD

- Moments of PDF given by local twist-2 operators (twist = dim - spin); limited to first few moments but carried out successfully

$$\langle x^n \rangle$$

Beyond the first few moments

- Smeared sources: Davoudi & Savage
- Gradient flow: Monahan & Orginos
- Current-current correlators: K.-F. Liu & S.-J. Dong; Braun & Müller; Detmold & Lin; QCDSF; Qiu & Ma
- Xiangdong Ji (Phys. Rev. Lett. 110 (2013) 262002): quasi-PDF: computing the x-dependence directly. (variation: pseudo-PDF, Radyushkin; w/ Karpie, Orginos, Zafeiropoulos)

Ji's observation

For a twist-2 operator,

$$\hat{\mathcal{O}}^{\mu\nu} = \bar{\psi} \left(\frac{i\gamma^\mu D^\nu + i\gamma^\nu D^\mu}{2} - \frac{i}{4} g^{\mu\nu} \not{D} \right) \psi$$

$$\langle P | \hat{\mathcal{O}}^{zz} | P \rangle = 2\langle x \rangle_q (P^z P^z - \frac{1}{4} g^{zz} M^2)$$

$$\langle P | \bar{\psi} i\gamma^z D^z \psi | P \rangle = 2\langle x \rangle_q P^z P^z \left[1 + \mathcal{O}\left(\frac{M^2}{P_z^2}\right) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{P_z^2}\right) \right]$$

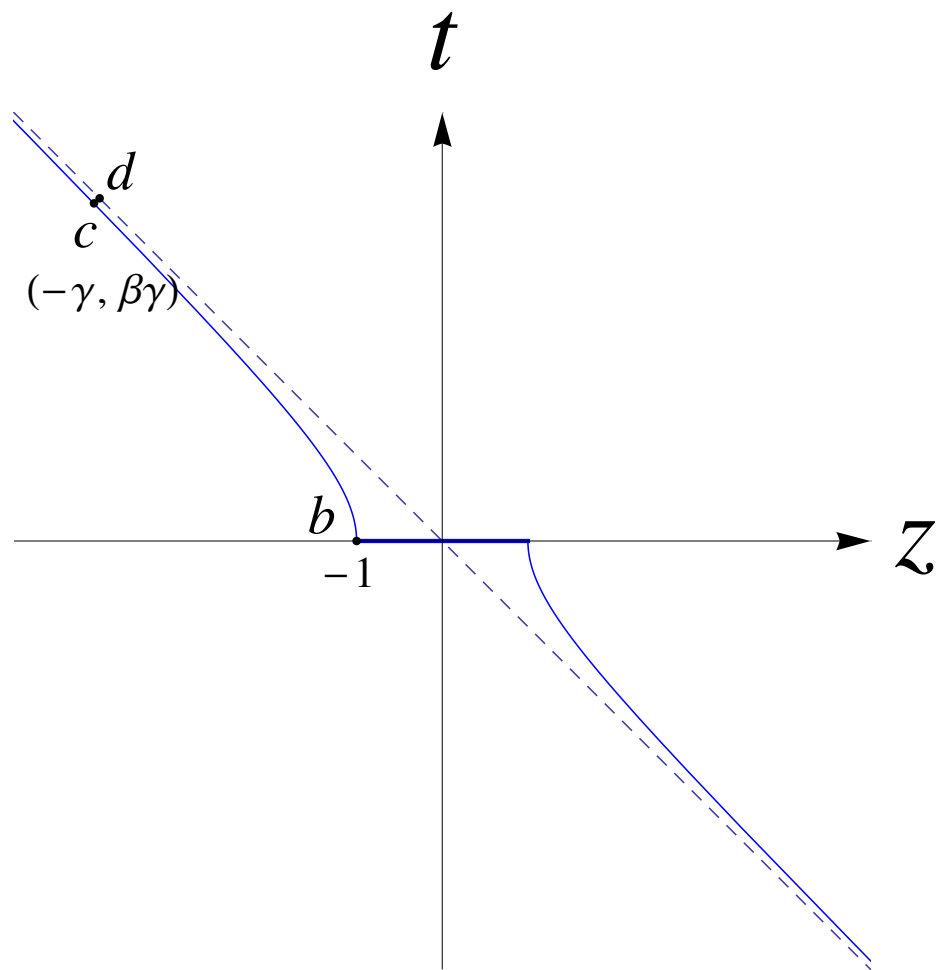
An equal time quark bilinear operator along the z direction gives the PDF of a proton moving in the z direction?

Ji's idea

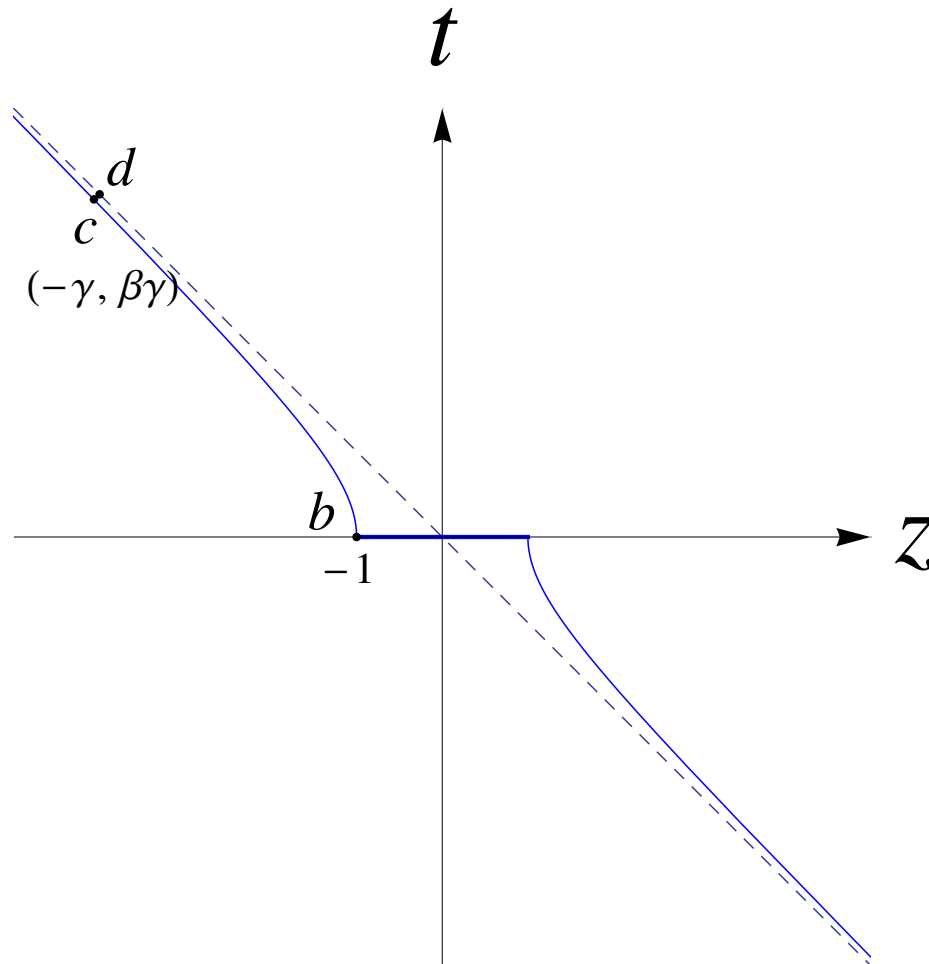
- Quark PDF in a proton: $(\lambda^2 = 0)$

$$q(x, \mu^2) = \int \frac{d\xi^-}{4\pi} e^{ix\xi^- P^+} \langle P | \bar{\psi}(0) \lambda \cdot \gamma \Gamma \psi(\xi^- \lambda) | P \rangle$$

- Boost invariant in the z-direction, rest frame OK
- Quark bilinear op. always on the light cone
- What if the quark bilinear is slightly away from the light cone (space-like) in the proton rest frame?



- Then one can find a frame where the quark bilinear is of equal time but the proton is moving.



- Analogous to HQET: need power corrections & matching---LaMET (Large Momentum Effective Theory)

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$$\tilde{q}(x, \mu^2, P^z) = \int \frac{dz}{4\pi} e^{-ixzP^z} \langle P | \bar{\psi}(0) \lambda \cdot \gamma \Gamma \psi(z\lambda) | P \rangle$$

$$\lambda^\mu = (0, 0, 0, 1)$$

$$\tilde{q}(x, \Lambda, P_z) = \int \frac{dy}{|y|} Z\left(\frac{x}{y}, \frac{\mu}{P_z}, \frac{\Lambda}{P_z}\right) q(y, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}^2}{P_z^2}, \frac{M^2}{P_z^2}\right) + \dots$$

UV, PQCD

Comment: Boosting makes the short distance expansion look bigger---physics depends on $z P_z$!

Power Corrections

JWC, Cohen, Ji, Lin, Zhang, hep-ph/1603.06664

M^2/P_z^2 corrections to all orders

$$\lambda_{\mu_1} \cdots \lambda_{\mu_n} P^{(\mu_1} \cdots P^{\mu_n)} = \lambda_{(\mu_1} \cdots \lambda_{\mu_n)} P^{\mu_1} \cdots P^{\mu_n}$$

(...) means the indices enclosed are symmetric and traceless

$$\lambda_{(\mu_1} \cdots \lambda_{\mu_n)} = \sum_{i=0}^{i_{\max}} B_{n,i} (\lambda^2)^i \left(\frac{\partial^2}{\partial \lambda_\alpha \partial \lambda^\alpha} \right)^i \lambda_{\mu_1} \cdots \lambda_{\mu_n}$$

$$i_{\max} = \frac{n - \text{Mod}[n, 2]}{2} \text{ and } B_{n,0} = 1$$

$$g^{\mu_1 \mu_2} P^{\mu_3} \cdots P^{\mu_n} \lambda_{(\mu_1} \cdots \lambda_{\mu_n)} = 0,$$

$$\sum_{i=0}^{i_{\max}} B_{n,i} (\lambda^2)^i \left(\frac{\partial^2}{\partial \lambda_\alpha \partial \lambda^\alpha} \right)^i \lambda^2 (\lambda \cdot P)^{n-2} = 0.$$

M^2/P_z^2 corrections to all orders

$$B_{n,i} = -\frac{B_{n,i-1}}{4i(n-i+1)}$$

$$(P^z P^z - \frac{1}{4}g^{zz}M^2)$$

.

M^2/P_z^2 corrections to all orders

$$K_n \equiv \frac{\langle x^{n-1} \rangle_{\tilde{q}}}{\langle x^{n-1} \rangle_q} = \frac{\lambda_{(\mu_1} \cdots \lambda_{\mu_n)} P^{\mu_1} \cdots P^{\mu_n}}{\lambda_{\mu_1} \cdots \lambda_{\mu_n} P^{\mu_1} \cdots P^{\mu_n}}$$

$$= \sum_{i=0}^{i_{\max}} C_{n-i}^i c^i,$$

where C is the binomial function and $c = -\lambda^2 M^2/4 (\lambda \cdot P)^2 = M^2/4P_z^2$ with $\lambda^\mu = (0, 0, 0, -1)$ and $\lambda \cdot P = P_z$.

M^2/P_z^2 corrections to all orders

$$\begin{aligned} q(x) &= \sqrt{1+4c} \sum_{n=0}^{\infty} \frac{f_-^n}{f_+^{n+1}} \left[(1+(-1)^n) \tilde{q}\left(\frac{f_+^{n+1}x}{2f_-^n}\right) + (1-(-1)^n) \tilde{q}\left(\frac{-f_+^{n+1}x}{2f_-^n}\right) \right] \\ &= \sqrt{1+4c} \sum_{n=0}^{\infty} \frac{(4c)^n}{f_+^{2n+1}} \left[(1+(-1)^n) \tilde{q}\left(\frac{f_+^{2n+1}x}{2(4c)^n}\right) + (1-(-1)^n) \tilde{q}\left(\frac{-f_+^{2n+1}x}{2(4c)^n}\right) \right] \end{aligned}$$

$$f_{\pm} = \sqrt{1+4c} \pm 1$$

$\Lambda_{\text{QCD}}^2 / P_z^2$ correction

$$\tilde{q}(x, \Lambda, P_z) \rightarrow \tilde{q}(x, \Lambda, P_z) + \tilde{q}_{\text{twist-4}}(x, \Lambda, P_z)$$

$$\tilde{q}_{\text{twist-4}}(x, \Lambda, P_z) = \frac{1}{8\pi} \int_{-\infty}^{\infty} dz \Gamma_0(-ixzP_z) \langle P | \mathcal{O}_{\text{tr}}(z) | P \rangle$$

$$\begin{aligned} \mathcal{O}_{\text{tr}}(z) = & \int_0^z dz_1 \bar{\psi}(0) \left[\gamma^\nu \Gamma(0, z_1) D_\nu \Gamma(z_1, z) \right. \\ & \left. + \int_0^{z_1} dz_2 \lambda \cdot \gamma \Gamma(0, z_2) D^\nu \Gamma(z_2, z_1) D_\nu \Gamma(z_1, z) \right] \psi(z\lambda). \end{aligned}$$

Γ_0 is the incomplete Gamma function $\int_0^1 \frac{dt}{t} e^{ix/t} = \Gamma_0(-ix)$

Power divergent mixing with \tilde{q} on the lattice?

Next Lecture: Matching Coefficients