

# **$B \rightarrow D^* \ell \nu$ decays on the lattice**

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**for JLQCD Collaboration**

seminar @ RIKEN R-CCS, February 17, 2021

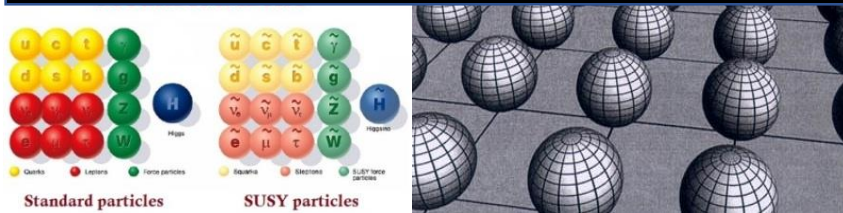
# principles and history of universe

super-symmetry?

Planck scale :  $O(10^{19})$  GeV

new physics

SUSY?, extra-dimension?, ...

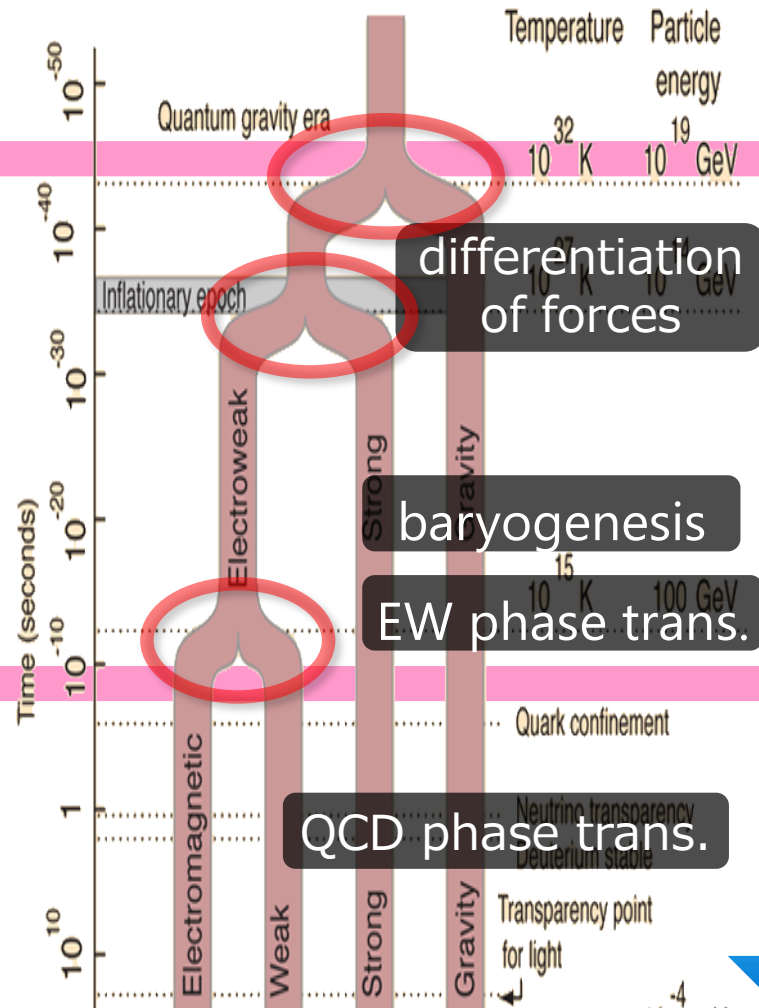


electro-weak scale :  $O(100)$  GeV

Standard Model

search for "new physics" beyond SM

inflation • big ban

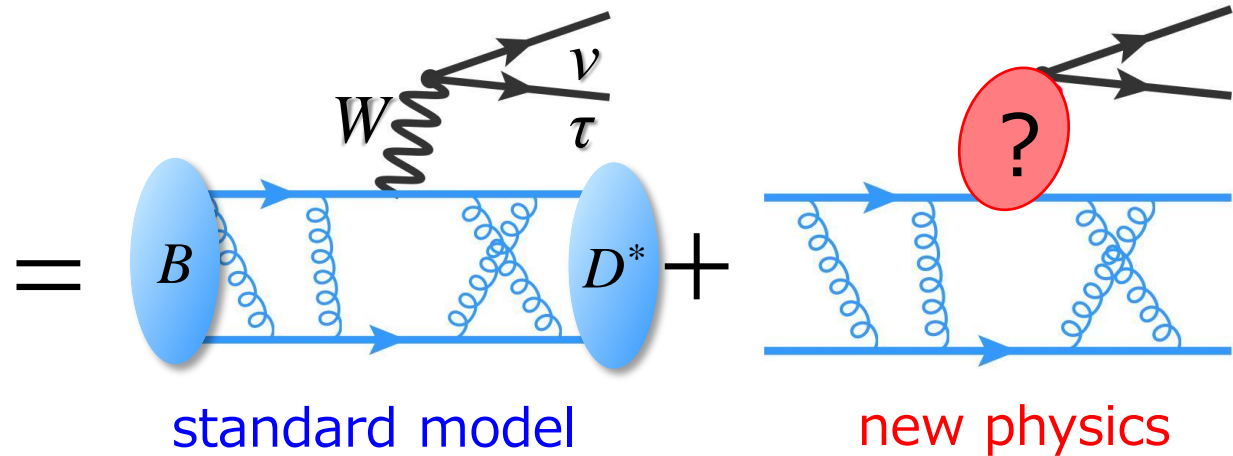
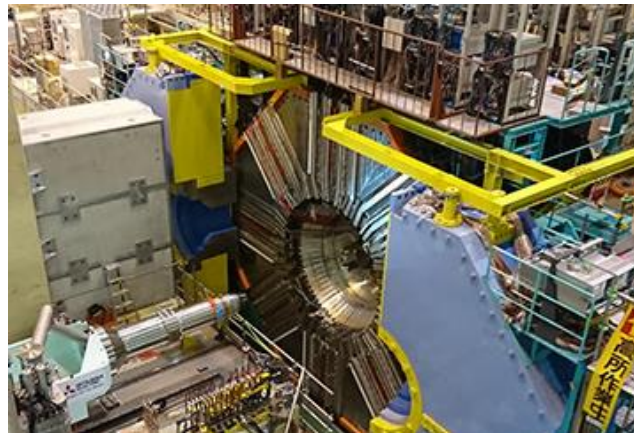


history of universe

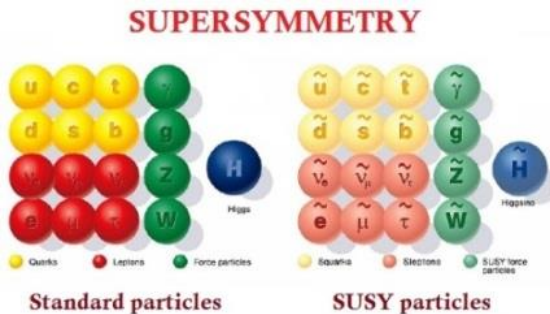
high energy

# intensity frontier for new physics

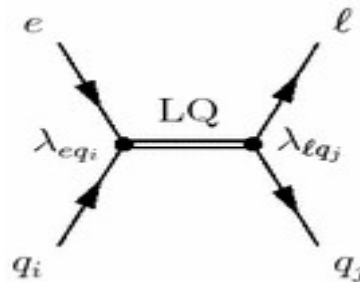
experiments vs standard model  $\Rightarrow$  new physics



charged Higgs ?



lepto-quark ?



- Higgs sector  $\Rightarrow$  EW symmetry breaking
- grand unification
- ...

# Outline

form factors of  $B \rightarrow D^* \ell \nu$  decays from lattice QCD

- $B \rightarrow D^* \ell \nu$  decays and form factors
  - background: flavor physics, “CKM” element  $|V_{cb}|$ , new physics
  - form factors: definition, phenomenology
- JLQCD 's simulation
  - simulation setup ... extrapolation to the real world
- for determination of  $|V_{cb}|$ 
  - recent determinations
  - indication of our lattice results

please ask questions any time

# 1. $B \rightarrow D^* \ell \nu$ decays and form factors

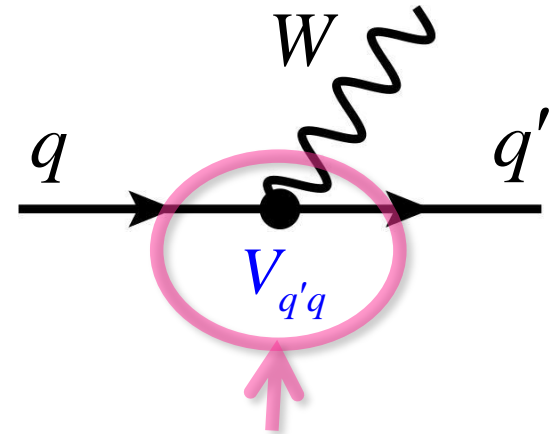
# flavor physics and CKM matrix

## flavor physics

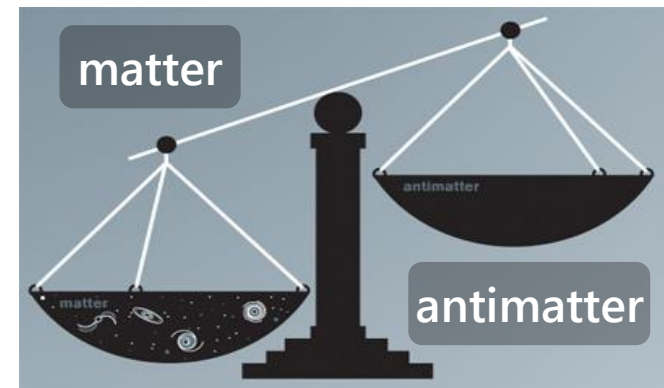
- processes w/ quark flavor change(s)
- Cabbibo-Kobayashi-Maskawa mechanism
- elements of 3x3 CKM matrix
  - ⇒ strength of flavor changes
  - ⇒ **amplitude of processes**

## missing anti-matter

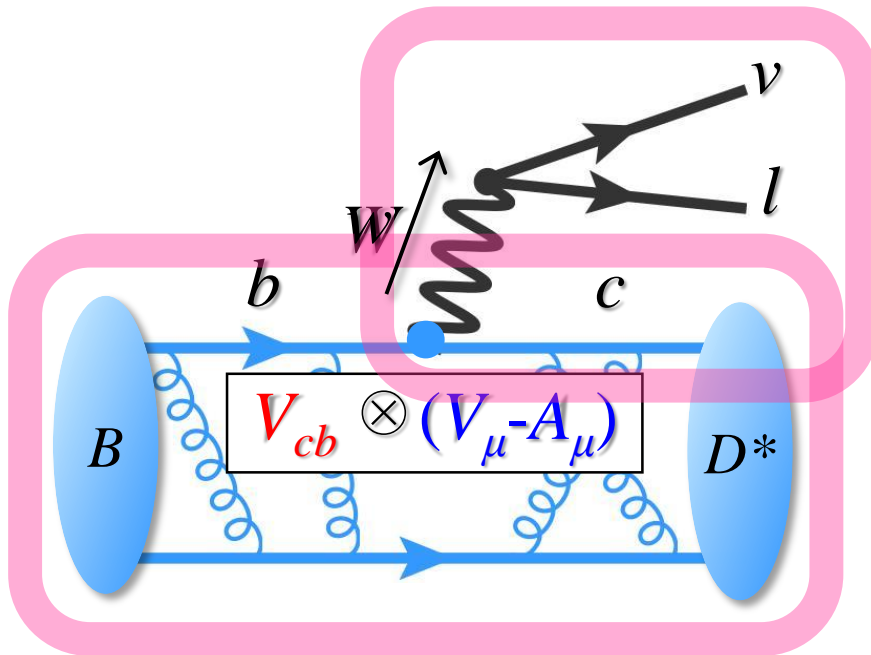
- a condition : violation of "CP" symmetry
  - C : particle  $\leftrightarrow$  anti-particle; P : parity trans.
- in SM : CKM and QCD  $\theta$ -term nEDM
- CKM is not sufficient to explain  $\rightarrow$
- **other sources (new physics) ???**



$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



# $B \rightarrow D^* \ell \nu$ semileptonic decay



- mediated by weak interaction  
 $b \rightarrow c W(\rightarrow \ell \nu) \Rightarrow$  pick up  $|V_{cb}|$
- $W$ +lepton part : perturbation
- hadronic QCD effects  $\Rightarrow$  lattice
- recoil :  $(\nu - \nu')^2 = 2 + w$ ,  $w = \nu \nu'$

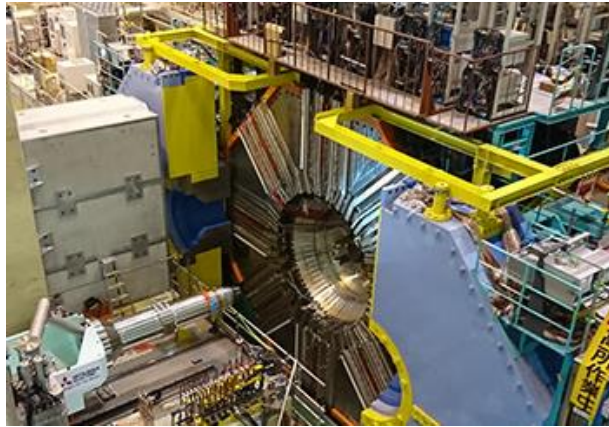
$$\mathcal{M}(B \rightarrow D^* \ell \nu) = V_{cb} \langle \ell \nu | V_\mu - A_\mu | 0 \rangle G_W \langle D^* | V_\mu - A_\mu | B \rangle$$

experiment
theory

- lepton+neutrino+W+hadron parts  $\Rightarrow$  determination of  $|V_{cb}|$
- $|V_{cb}|$  from another decay, ratio  $\Rightarrow$  search for new physics



# experiments



KEKB / Belle  
@KEK (-2010)

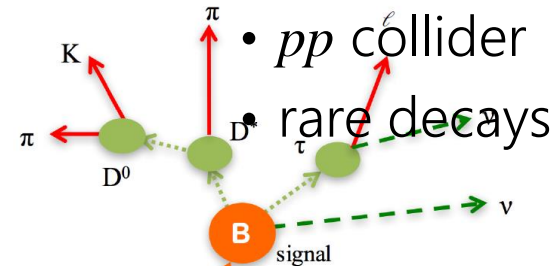


BaBar @ SLAC (-2008)

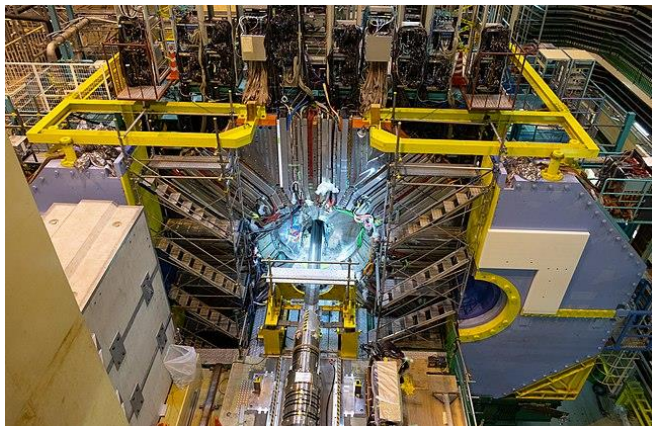
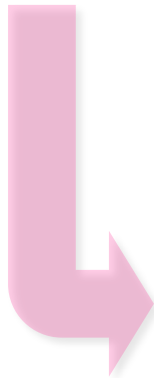


LHCb @ CERN  
Run3 2023-

- $e^+e^- \rightarrow O(10^9)$  pairs of  $BB, B^+B^-$
- good at semileptonic decays w/  $\nu$



- $pp$  collider
- rare decays



SuperKEKB / Belle II@KEK (2019-)

- upgrade of Belle
- $\times 50$  statistics by 2031
- "tagged" data
- "untagged" is also possible

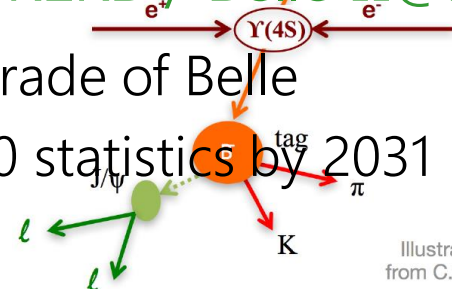
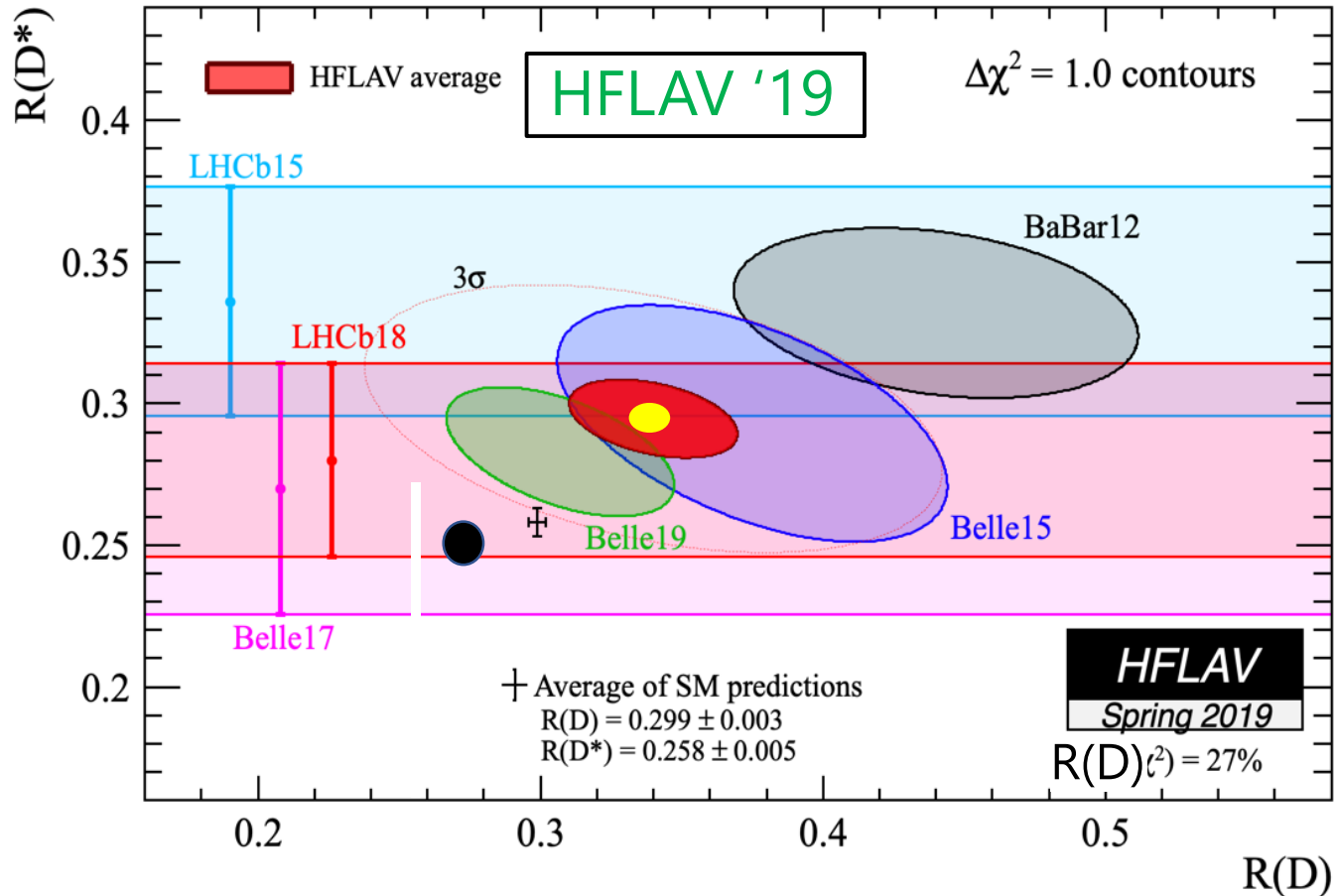


Illustration  
from C. Bozzi

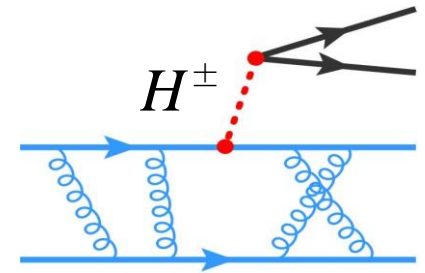


# promising probe of new physics

$$R(D^{(*)}) = \Gamma(B \rightarrow D^{(*)} \tau \nu) / \Gamma(B \rightarrow D^{(*)} \{e, \mu\} \nu)$$



- 3.4  $\sigma$  from SM
- $\tau$  channel via  $b \rightarrow c\tau\nu$   
 $\Rightarrow$  charged Higgs?

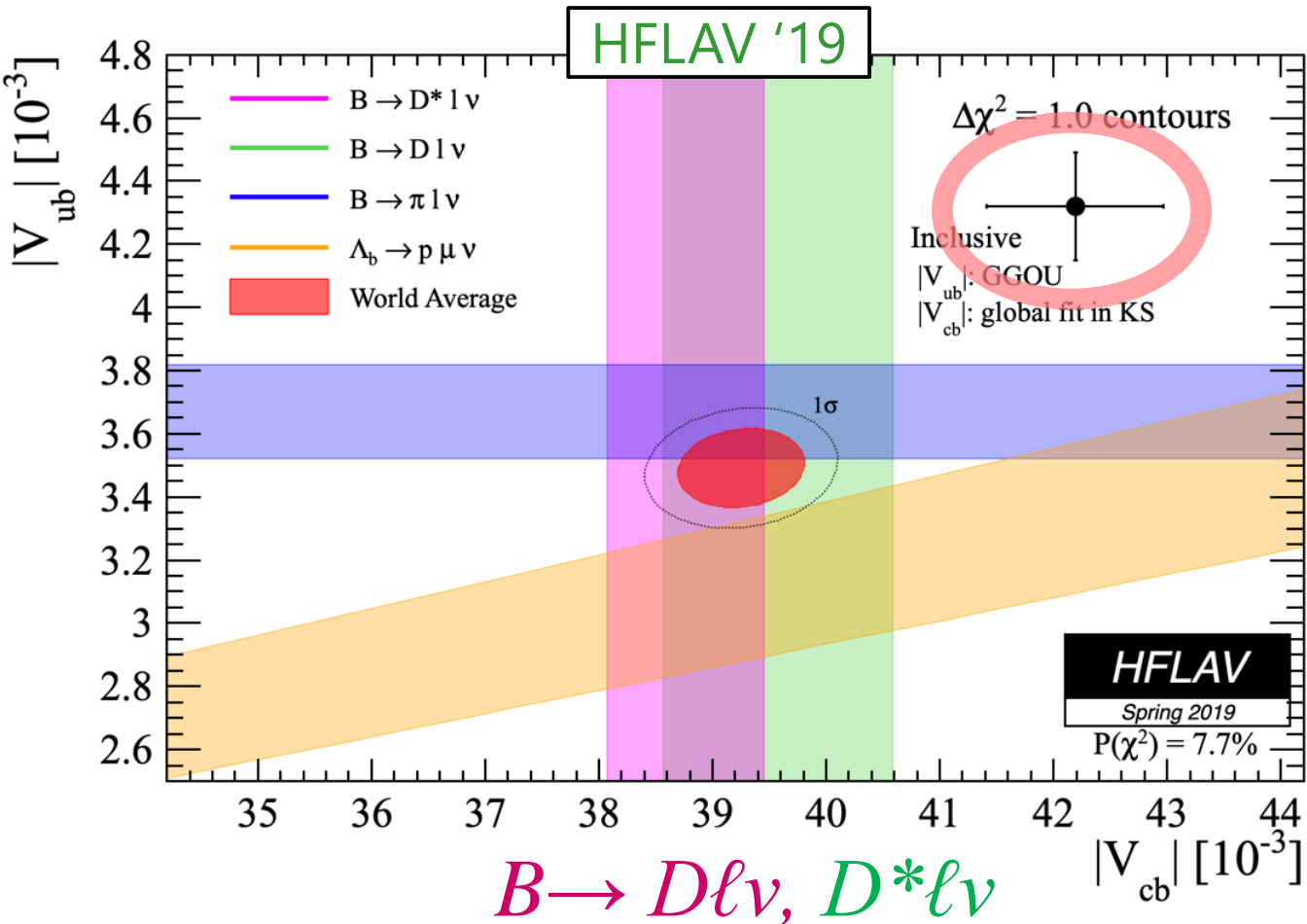


- inconsistent exp's?

Belle II : 3  $\sigma$  tension  $\Rightarrow$  10  $\sigma$  evidence !?

# determination of CKM element $|V_{cb}|$

long-standing discrepancy (6-15%)



inclusive decay  
= sum up over final  
state hadron(s)

$$B \rightarrow X_u \ell \nu$$

$$= \pi \ell \nu + \rho \ell \nu + \dots$$

exclusive decay

$$B \rightarrow \pi \ell \nu$$

to be resolved for precision new physics search

# hadronic matrix elements

“relativistic” convention

$$\langle D^*(p', \epsilon) | V^\mu | \bar{B}(p) \rangle = ig \epsilon^{\mu\alpha\beta\gamma} \epsilon_\alpha^* p'_\beta p_\gamma,$$

$$\langle D^*(p', \epsilon) | A^\mu | \bar{B}(p) \rangle = f \epsilon^{*\mu} + (\epsilon^* \cdot p) [a_+ (p + p')^\mu + a_- (p - p')^\mu]$$

“heavy quark” convention ( $\sqrt{M_B} |B\rangle \rightarrow |B\rangle$ )

$$\langle D^*(v', \epsilon') | V^\mu | \bar{B}(v) \rangle = ih_V(w) \epsilon^{\mu\nu\alpha\beta} \epsilon_\nu'^* v'_\alpha v_\beta$$

$$\langle D^*(v', \epsilon') | A^\mu | \bar{B}(v) \rangle = h_{A_1}(w) (w + 1) \epsilon'^{* \mu} - [h_{A_2}(w) v^\mu + h_{A_3}(w) v'^\mu] \epsilon^* \cdot v$$

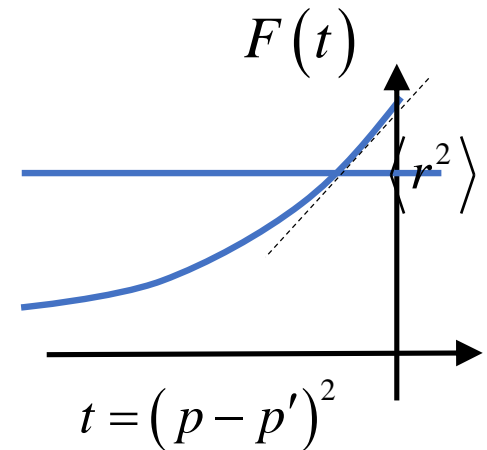
- 4 form factors (FFs) w/  $\epsilon_{D^*} p_{D^*} = 0$
- function of momentum transfer / recoil parameter
- $q^2 = t = (p - p')^2$ ,  $q^2 \leq (M_B - M_{D^*})^2 = q^2_{\max} = t_-$  ( $\mathbf{p} = \mathbf{p}' = \mathbf{0}$ )
- $v = p/M_B$ ,  $w = v \cdot v' = (M_B^2 + M_{D^*}^2 - q^2)/2M_B M_{D^*}$ ,  $1 \leq w$

# "form" factors (FFs)

Lorentz invariant function of Lorentz invariant variable describing MEs

*cf.* pion electro-magnetic form factor

- form factor of  $\pi \rightarrow \pi\gamma$
- shape w.r.t charge distribution probed by  $\gamma$ 
  - point-like particle  $\Rightarrow$   $t$ -indep. constant
  - slope = charge radius



$B \rightarrow D^* \ell \nu$  form factors

- shape of B probed by W boson
- shape of QCD bound state  $\Rightarrow$  essentially non-perturbative

**how can we describe FFs analytically?**

# recoil parameter dependence

small parameter expansions

zero recoil expansion in  $(w - 1)$

- heavy quark convention, near  $w = 1$ , interpolation

z-parameter expansion

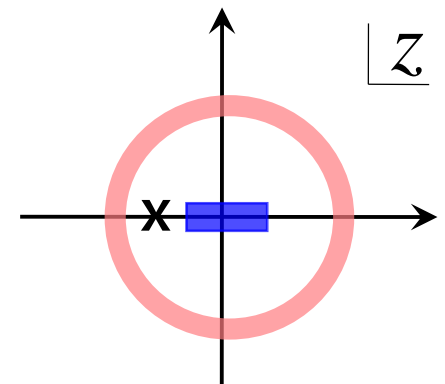
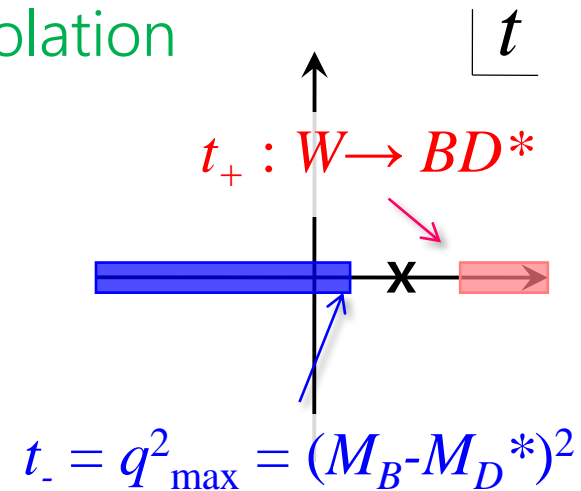
- map semileptonic region to short segment
- factor out singularity due to resonances



$$z = \frac{\sqrt{t_+ - t} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - t} + \sqrt{t_+ - t_0}}$$

$$F(t) = \frac{1}{P(t)\phi(t; t_0)} \sum_{n=0}^{\infty} a_n z(t; t_0)^n \quad P(t) = t - M_{\text{reso}}^2$$

- model-independent  $\Rightarrow$  phenomenology





# quark mass dependences

## heavy quark effective theory (HQET)

- based on heavy quark symmetry w.r.t heavy hadron's flavor and spin
- FFs:  $1/m_Q^n$  expansion around the  $m_Q = \infty$  limit
- NLO is known but with a set of unknown parameters (Isgur-Wise func.)
- unknown higher order corrections MAY be small for FF ratios  $h_X / h_Y$

## chiral perturbation theory (ChPT)

- based on spontaneous breaking of chiral symmetry w.r.t. light flavors
- non-analytic correction "chiral logarithm"  $\ln[M_\pi^2/\mu^2]$
- NLO but with a set of unknown parameters (LECs)
- heavy quark symmetry  $\Rightarrow$  suppressed by  $D^*-D$  mass splitting

w/ unknown param.s but useful in LQCD

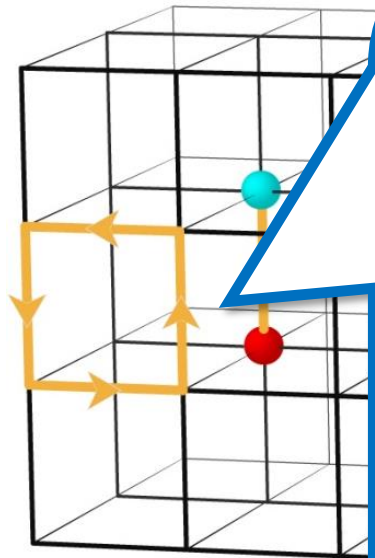
## **2. JLQCD collaboration's study**

# lattice QCD

numerical and first-principles calculation



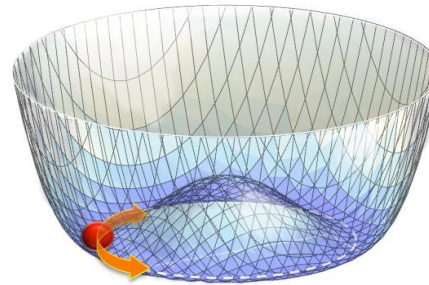
QCD on a space-time



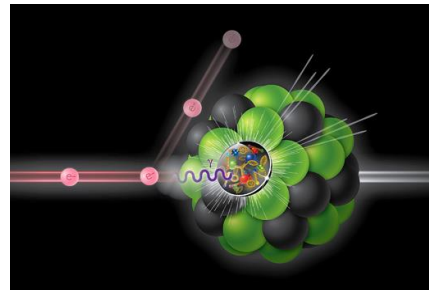
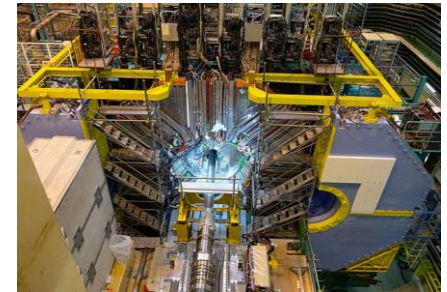
massively parallel  
e.g. on Oakforest  
HPCI: hp170106,

wide applications

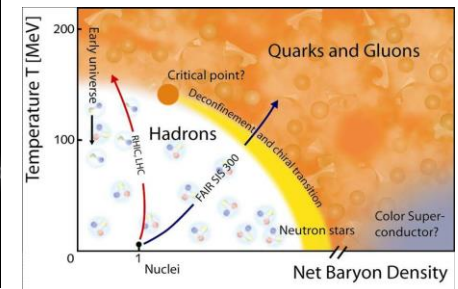
nature of QCD



new physics search



nuclear physics



phase structure

# JLQCD's simulation of QCD ①

preserving chiral symmetry

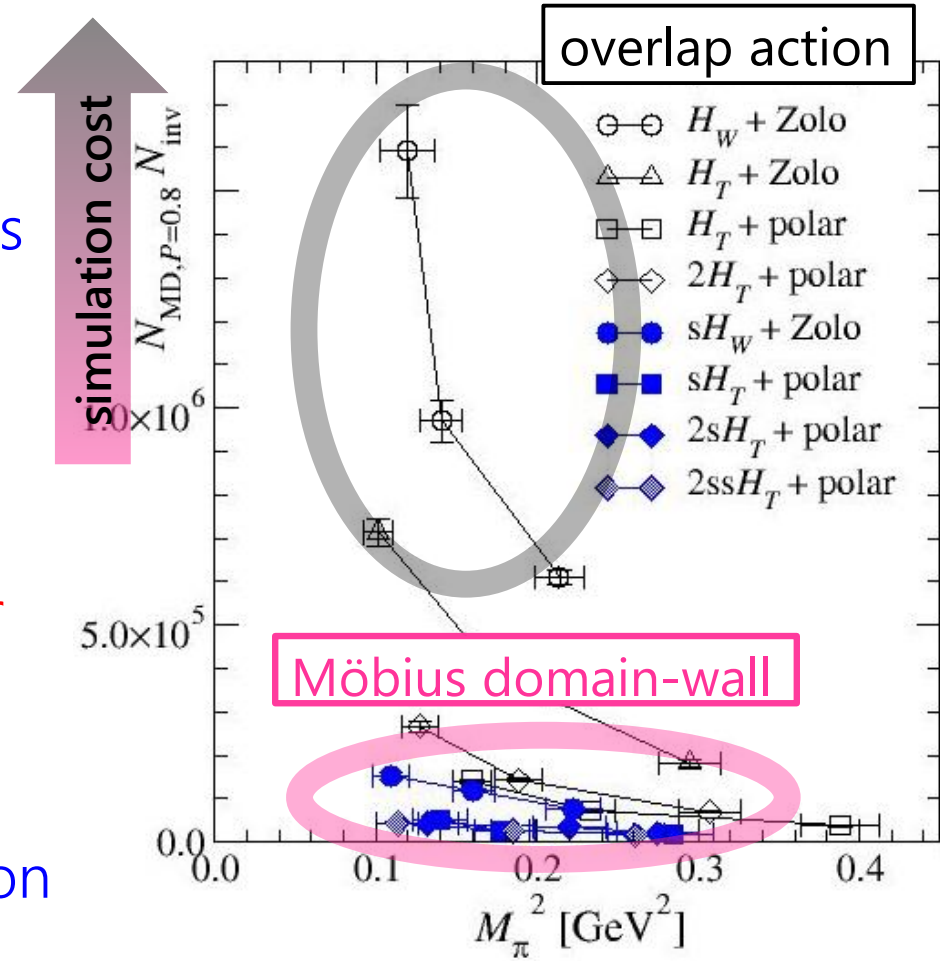
## chiral symmetry

- spontaneous breaking:  
characterize low-energy dynamics
  - ⇔ Wilson types : explicitly violated
  - ⇔ staggered types : partly, locality

## for heavy quark physics

- forbid leading discretization error
- simplified renormalization
  - lattice : removed  $\infty$ 's
  - lattice ⇔  $\overline{\text{MS}}$  : finite correction
- ChPT to describe  $m_q$  dependence

computationally demanding ⇒ speed up x 20

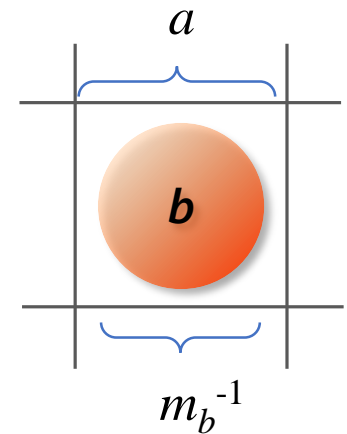
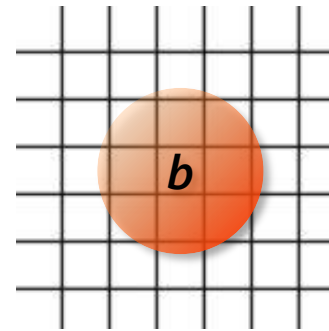


# JLQCD's simulation of QCD ②

## relativistic QCD

### relativistic QCD

- QCD as it is
  - lattice fine enough for  $m_b$ ?
    - Compton wavelength  $1/m_b$
    - discretization error as  $O((am_b)^n)$
- ⇒ unphysically small  $m_b$  + extrapolation to  $m_{b,\text{phys}}$



### effective theory (ET) approach

- NRQCD / HQET action for heavy quarks
- can simulate  $m_{b,\text{phys}}$
- need matching : parameter tuning of ET to reproduce QCD

independent calculations w/ different systematics



# previous and on-going studies

only limited number of studies

## published studies

- Fermilab / MILC '14    staggered light quarks + ET-based heavy quarks
- HPQCD '17                staggered light quarks + relativistic heavy quarks
- only  $h_{A1}$  @  $w=1$      $\Leftrightarrow$  four FFs  $h_{A1}, h_{A2}, h_{A3}, h_V$

## need to calculate all FFs

## on-going studies

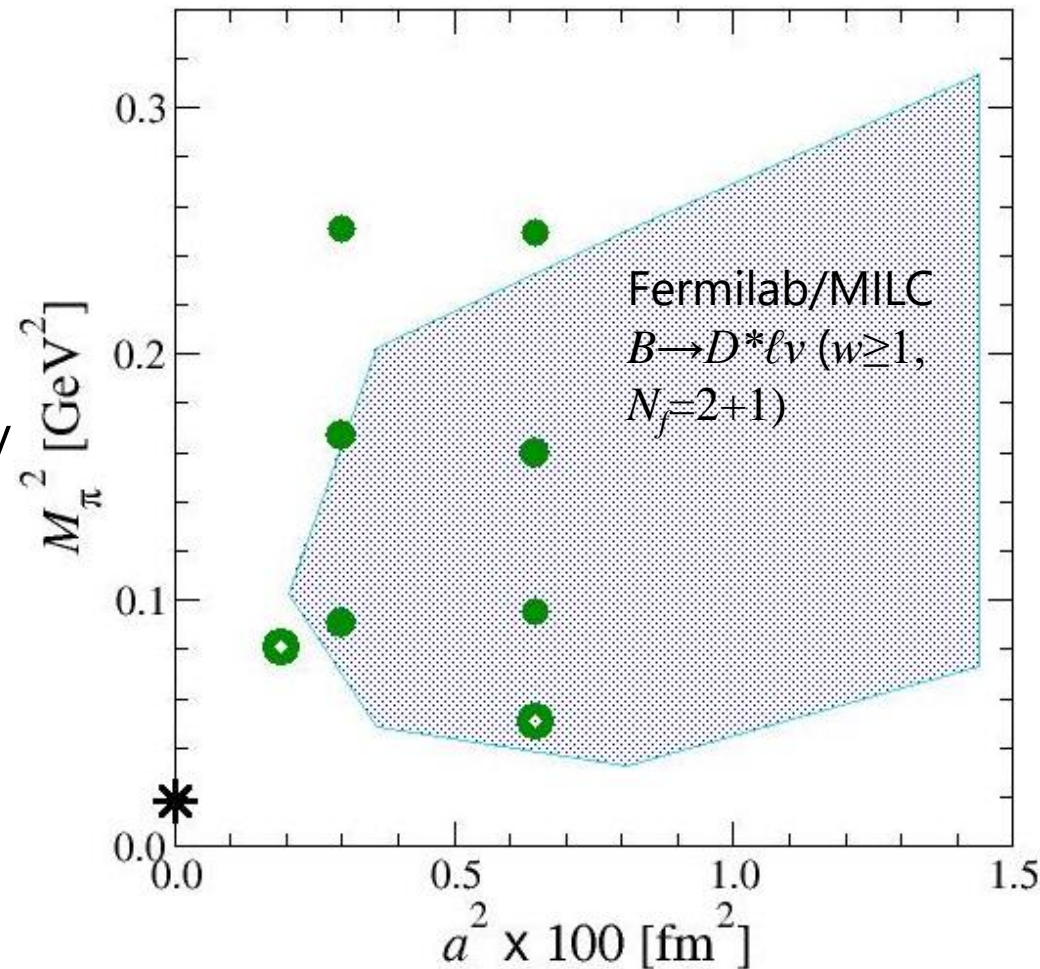
- Fermilab / MILC                staggered light quarks + ET-based heavy quarks
- JLQCD                                chiral light quarks                + relativistic heavy quarks
- calculating all FFs w/ very different systematics

# gauge ensembles

clean simulation w/ chiral symmetry

## simulation parameters

- $N_f = 2+1$
- $a^{-1} \sim 2.5, 3.6, 4.5 \text{ GeV}$
- $M_\pi \sim 230, 300, 400, 500 \text{ MeV}$
- 5,000 HMC traj. for each
- $M_\pi L \geq 4$



similar parameter regions w/ chiral symmetry

# measurements

## relativistic approach for hadron correlators

chiral fermion action  
also for heavy quarks

- $m_b = \{1.00, 1.25, 1.25^2, \dots\} \times m_c$

- $m_b \leq 0.7a^{-1}$

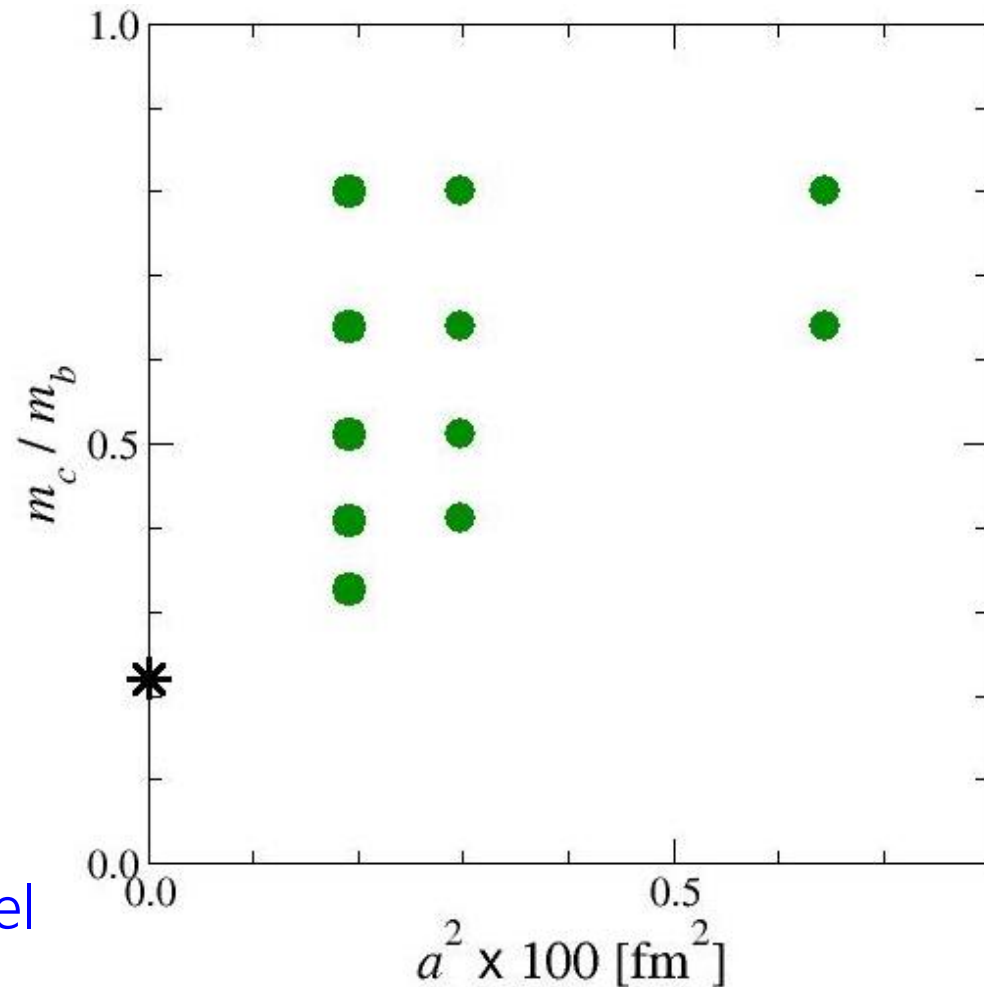
- extrapolation to  $m_{b,\text{phys}}$

⇔ no  $O(am_b)$  error

small  $O((am_b)^2)$  error

⇔ discretization effects

suppressed to a few % level



extrapolation is controllable

# ratio method (Hashimoto et al. '99)

$$\mathbf{p} = \mathbf{0}$$

$$|\mathbf{p}'|^2 = 0, 1, 2, 3, 4 \text{ in units of } (2\pi/L)^2, \text{ physical } m_c$$

$$\frac{\langle D^* | V_\mu^{(\text{lat})} | B \rangle}{\langle D^* | A_\mu^{(\text{lat})} | B \rangle} \rightarrow \frac{h_V(w)}{h_{A_1}(w)}$$

$$\langle B | O_B^\dagger \rangle, \exp[-M_B \Delta t], \dots \text{cancel}$$

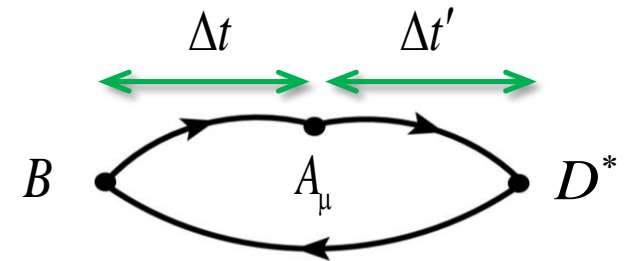
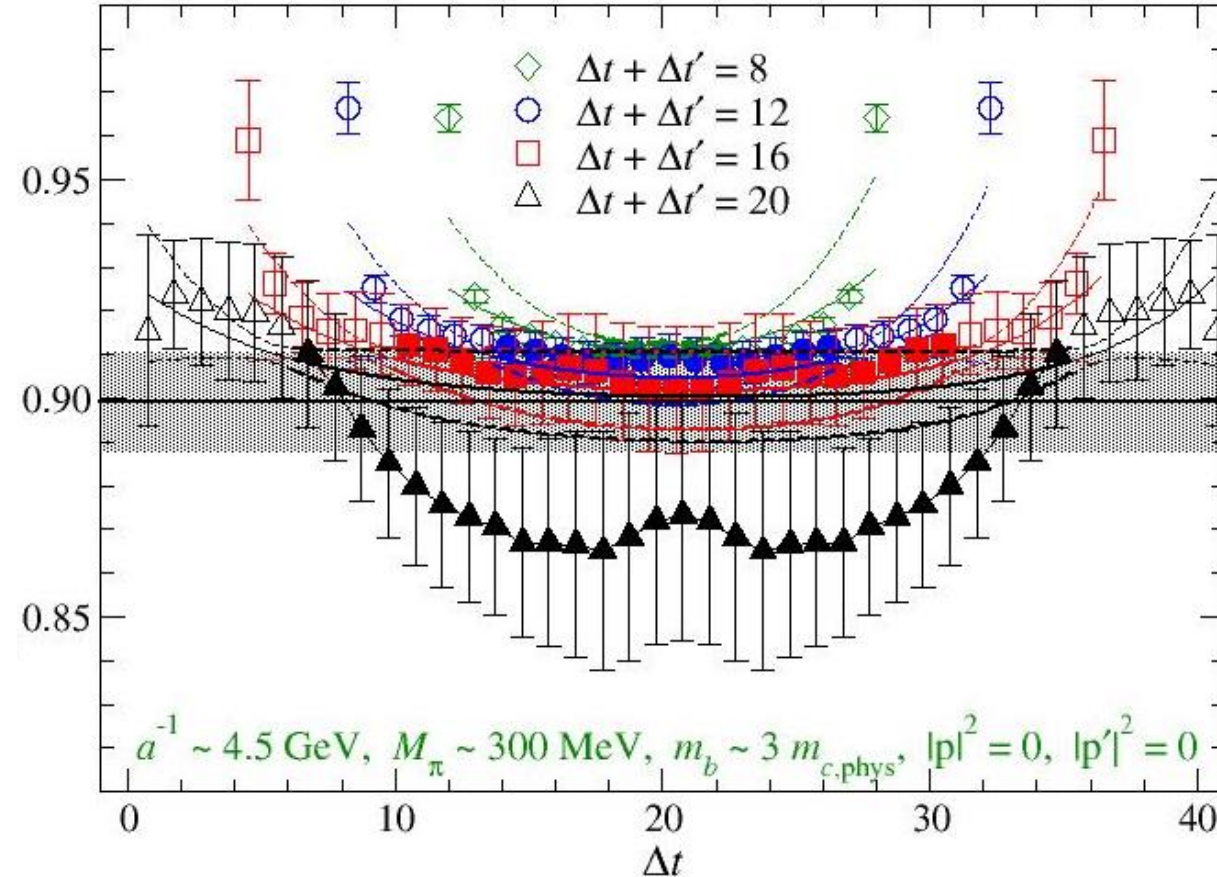
$$Z_A, Z_V \text{ cancel}$$

can calculate SM FFs w/o explicit renormalization

$$\{h_{A_2}(w), h_{A_3}(w), h_V(w)\} / h_{A_1}(w), \quad h_{A_1}(w) / h_{A_1}(1), \quad h_{A_1}(1) / \sqrt{F_B^{\text{EM}}(1) F_{D^*}^{\text{EM}}(1)}$$

# ground state contribution

$$\langle D^* | A_1 | B \rangle \langle B | A_1 | D^* \rangle / \langle D^* | V_4 | D^* \rangle \langle B | V_4 | B \rangle \rightarrow h_{A1} (1)$$



4 values of total separation  $\Delta t + \Delta t'$

$$\left| \langle D^* | A | B \rangle \right|^2 \times \left\{ 1 + c e^{-\Delta E_B \Delta t} + c' e^{-\Delta E_{D^*} \Delta t'} \right\}$$

multiple  $(\Delta t + \Delta t')$ 's  $\Rightarrow$  ~~excited state contamination~~  
good statistical accuracy



# extrapolation to the real world

**NLO HMChPT** (Randall-Wise '92, Savage '01)

- non-analytic log term (but small)
- polynomial in light quark masses

$$\xi_\pi = \frac{M_\pi^2}{(4\pi f_\pi)^2}, \quad \xi_{\eta_s} = \frac{M_{\eta_s}^2}{(4\pi f_\pi)^2}$$

$$h_{A1}(w) = c(m_c) + \frac{g_{D^*D\pi}^2}{16\pi^2 f_\pi^2} \Delta_c^2 b_{\log} \bar{F}_{\log}(M_\pi, \Delta_c, \Lambda_\chi) + c_\pi \xi_\pi + c_{\eta_s} \xi_{\eta_s}$$

$$+ c_w (w-1) + d_w (w-1)^2 + c_b \varepsilon_b + c_a \xi_a + c_{amb} \xi_{amb}$$

**0 recoil expansion**

- interpolation w/ quad.

**HQET**

- in  $\varepsilon_b = \frac{\bar{\Lambda}}{2m_b}$

**discretization effects**

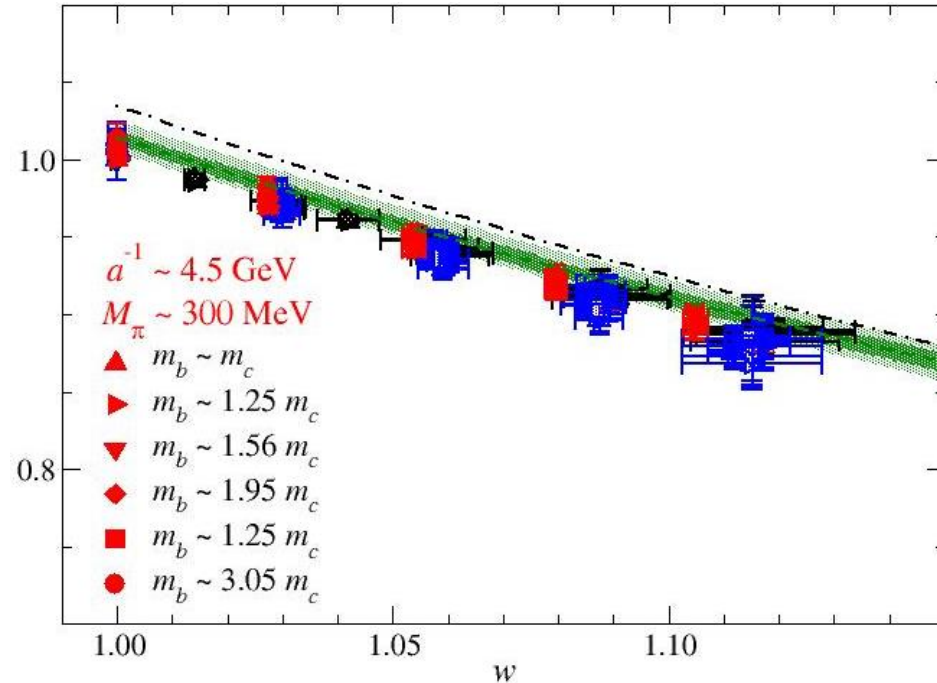
- in  $\xi_a = (a\Lambda_{\text{QCD}})^2$ ,  $\xi_a = (am_b)^2$

- systematic error of this form : adding higher order / remove terms
- large covariance matrix, input  $g_{D^*D\pi}$  , ...

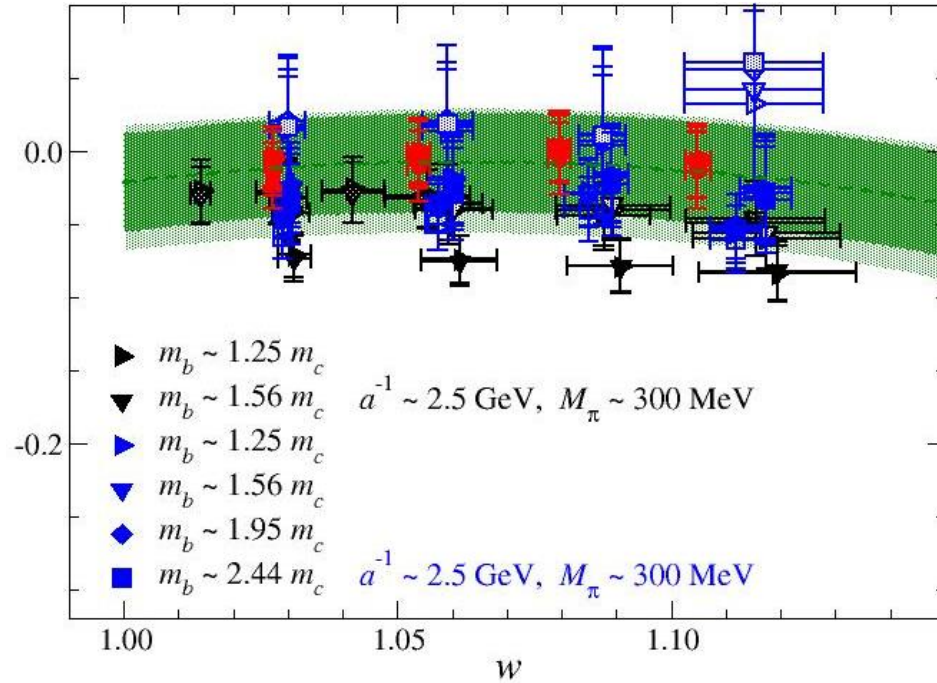
analysis is being finalized  $\Rightarrow$  preliminary results in this talk

# $B \rightarrow D \ell \nu$ form factors

$h_+$  VS  $w$



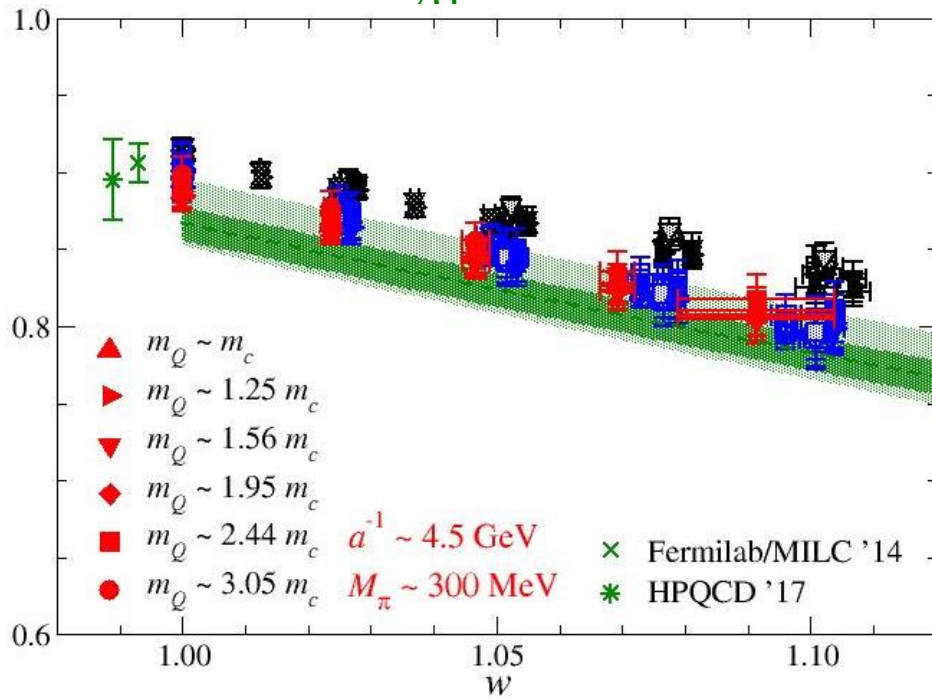
$h_-$  VS  $w$



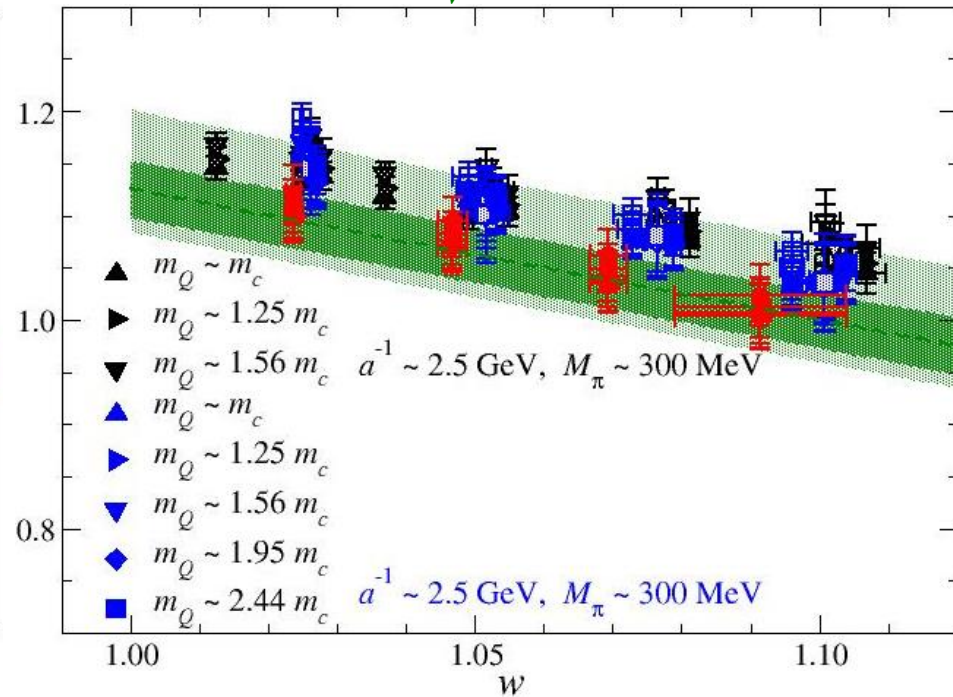
- mildly depend on  $a, M_\pi, m_s, m_b \Rightarrow \geq 50\%$  error except  $c, c_w, c_b$  ( $h_+$ )  
 $\Rightarrow$  consistent w/ world average (FLAG4 '19) within  $2\sigma$
- $h_+$ :  $\sim 1\%$  stat /  $1\%$  systematic errors @  $a = 0, m_{q,\text{phys}}$

# $B \rightarrow D^* \ell \nu$ form factors

$h_{A1}$  VS  $w$



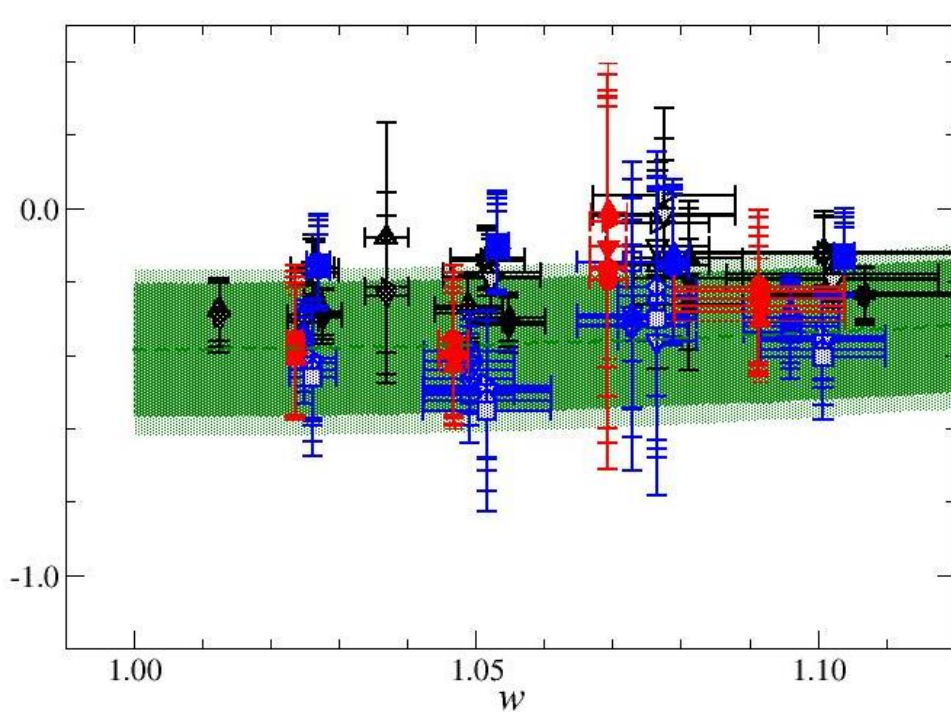
$h_V$  VS  $w$



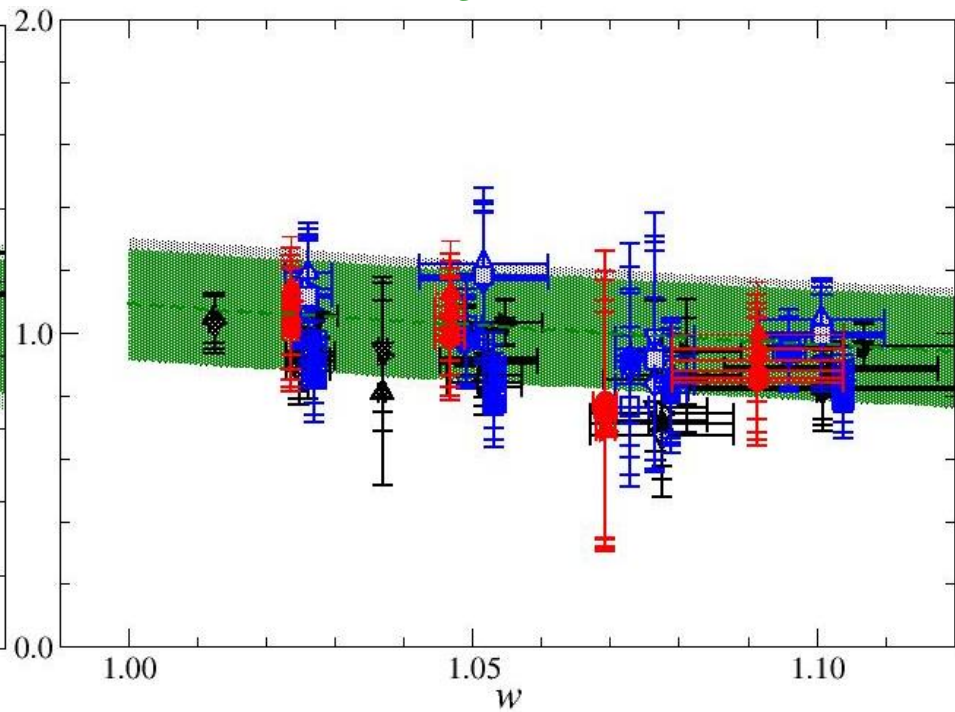
- mild  $a, M_\pi, m_s, m_b$  dependence
- only  $h_{A1}(1)$  from previous studies [consistent]
- $h_{A1}, h_V \sim 1 - 2 \% \text{ stat} / 3 - 6 \% \text{ systematic } (a \neq 0)$

# $B \rightarrow D^* \ell \nu$ form factors

$h_{A_2}$  VS  $w$



$h_{A_3}$  VS  $w$



- $h_{A_2}, h_{A_3} \sim 20\%$  stat. / sys. errors  $\Leftrightarrow \Delta[d\Gamma/dw] \sim$  Belle

$$\langle D^*(p', \varepsilon') | A_\mu | B(p) \rangle \Rightarrow \{h_{A_1}(w), h_{A_2}(w), h_{A_3}(w)\}$$

### **3. determination of $|V_{cb}|$**

# $B \rightarrow D^* \ell \nu$ differential decay rate

$B \rightarrow D^* \{e, \mu\} \nu$  in the limit  $m_l^2 = 0$

expr't

CKM matrix element

EW : perturbation

$$\frac{d\Gamma}{dw} = |V_{cb}|^2 \frac{G_F^2}{48\pi^3} (M_B - M_{D^*})^2 M_{D^*}^3 \sqrt{w^2 - 1} (w+1)^2 \eta_{EW} \times \left[ 2 \frac{1 - 2wr + r^2}{(1-r)^2} \left\{ 1 + \frac{w-1}{w+1} R_1(w) \right\} + \left\{ 1 + \frac{w-1}{1-r} (1 - R_2(w)) \right\}^2 \right] h_{A1}(w)^2$$

QCD : non-perturbative

- described by  $h_{A1}$ ,  $R_1 = \frac{h_V}{h_{A1}}$ ,  $R_2 = \frac{r h_{A2} + h_{A3}}{h_{A1}}$  (functions of  $w$ )

$|V_{cb}|$  : a relative normalization by comparing exp. and th.



# conventional determination of $|V_{cb}|$

$$\frac{d\Gamma}{dw} \propto |V_{cb}|^2 \left[ 2 \frac{1-2wr+r^2}{(1-r)^2} \left\{ 1 + \frac{w-1}{w+1} R_1(w) \right\} + \left\{ 1 + \frac{w-1}{1-r} (1-R_2(w)) \right\}^2 \right] h_{A1}(w)^2$$

- no published results for FFs except  $h_{A1}(1)$  !!
- ① assume a parametrization of FFs as a function of  $w$
- ② fix  $|V_{cb}|$  and unknown parameters by the fit to exp. and th. data  
 $\Rightarrow$  estimate FFs from experimental data !!

## Boyd-Grinstein-Lebed (BGL) parametrization '97

- model-independent  $z$  expansion based on analyticity
- many parameters ... can not be determined by previous exp. data

$$g(z) = \frac{1}{P_{1-}(z)\phi_g(z)} \sum_{n=0}^{\infty} a_n^g z^n \quad f(z) = \frac{1}{P_{1+}(z)\phi_f(z)} \sum_{n=0}^{\infty} a_n^f z^n \quad \mathcal{F}_1(z) = \dots$$

# conventional determination of $|V_{cb}|$

$$\frac{d\Gamma}{dw} \propto |V_{cb}|^2 \left[ 2 \frac{1-2wr+r^2}{(1-r)^2} \left\{ 1 + \frac{w-1}{w+1} R_1(w) \right\} + \left\{ 1 + \frac{w-1}{1-r} (1-R_2(w)) \right\}^2 \right] h_{A_1}(w)^2$$

## Caprini-Lellouch-Neubert (CLN) parametrization '97

- BGL + NLO HQET constraints  $\Rightarrow$  (too much?) less parameters
  - HQET for  $h_{A_1}(w) / h_{A_1}(1)$ ,  $R_1$ ,  $R_2$  expecting small correction to FF ratios

$$h_{A_1}(w) = h_{A_1}(1) \left( 1 - 8\rho_{D^*}^2 z + (53\rho_{D^*}^2 - 15) z^2 - (231\rho_{D^*}^2 - 91) z^3 \right)$$

fit parameters

NLO HQET

$$R_1(w) = R_1(1) - 0.12(w-1) + 0.05(w-1)^2$$

$$R_2(w) = R_2(1) + 0.11(w-1) - 0.06(w-1)^2$$

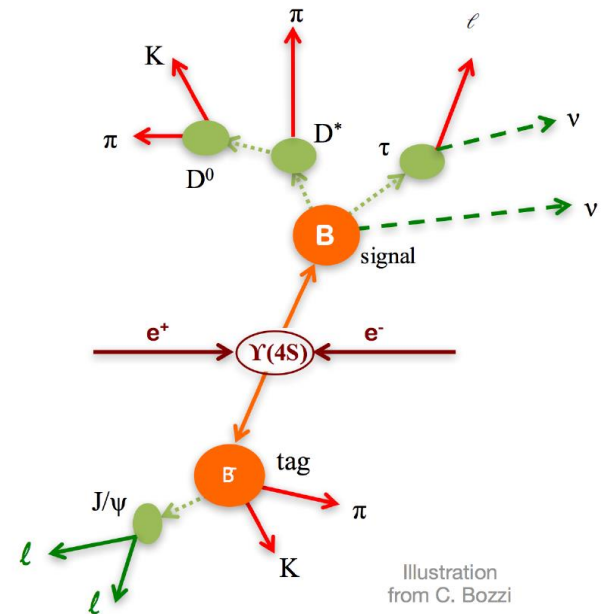
# determinations w/ Belle data

Belle '17 and '18 (near full statistics)

- full kinematical distribution ( $q^2$  and 3  $\theta$ 's)
  - '17: tagged (less miss-identification of  $\nu$ )
  - '18: untagged (more statistics)

conventional CLN analysis w/ tagged data

- Bernlochner - Ligeti - Papucci - Robinson '17



“model-independent” BGL analyses w/ tagged data

- Bigi – Gambino - Schacht '17 :  $|V_{cb}| \times 10^3 = 41.7 (+2.0/-2.2)$
- Grinstein – Kobach '17 :  $|V_{cb}| \times 10^3 = 41.9(+2.0/-1.9)$
- consistent w/ inclusive determination  $42.0 (0.5)$

HQET constraints for CLN are source of the  $|V_{cb}|$  tension?

# determinations w/ Belle data

it is plausible that a tension is resolved by removing phenomenological assumptions ...

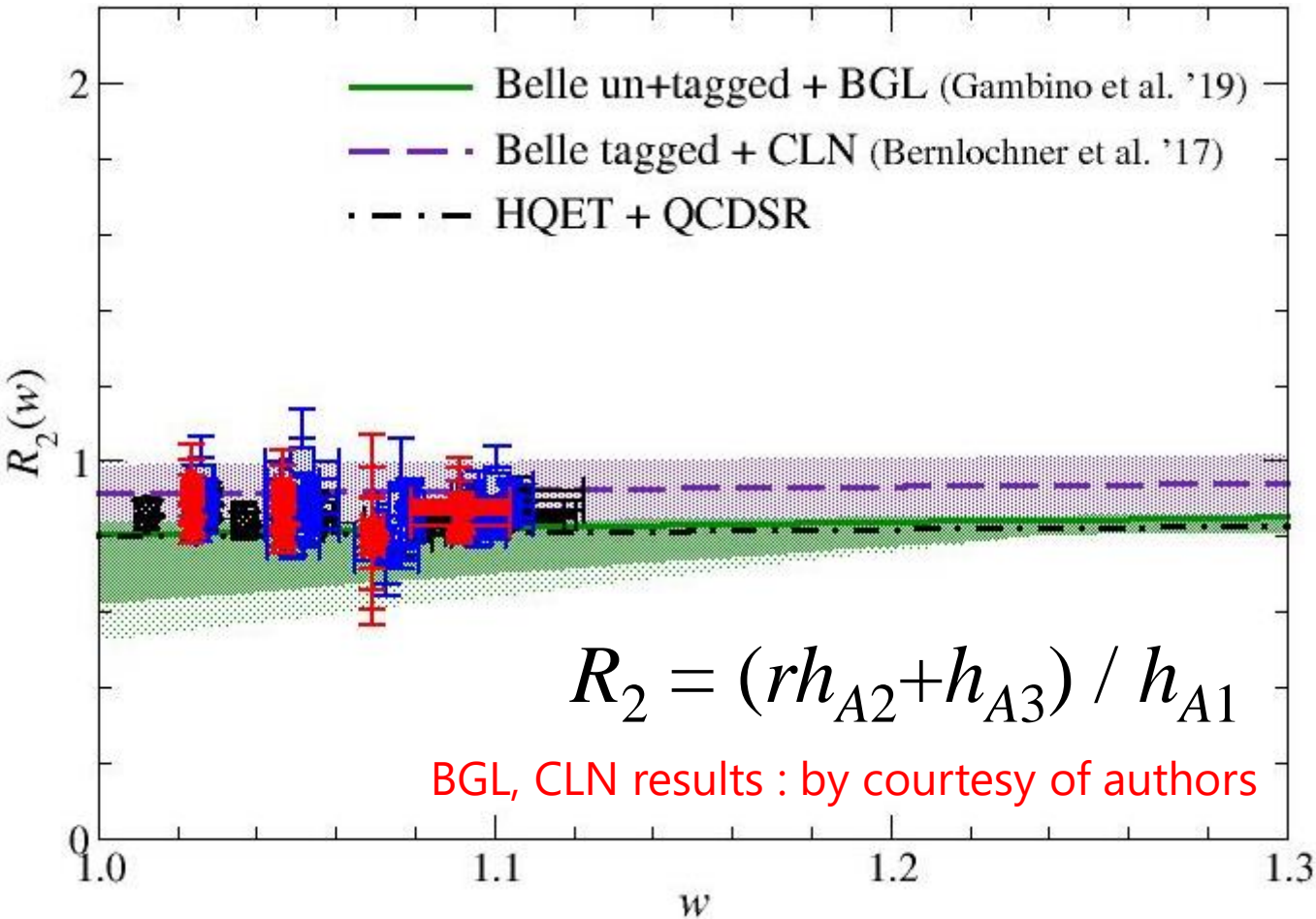
form factors should be predicted from QCD

the CLN and BGL analyses gave inconsistent results for

$$h_{A_1}(w) / h_{A_1}(1), R_1$$

⇒ let us compare w/ our lattice results from first-principles

# FF ratio $R_2$



no big difference  
among CLN, BGL,  
HQET...

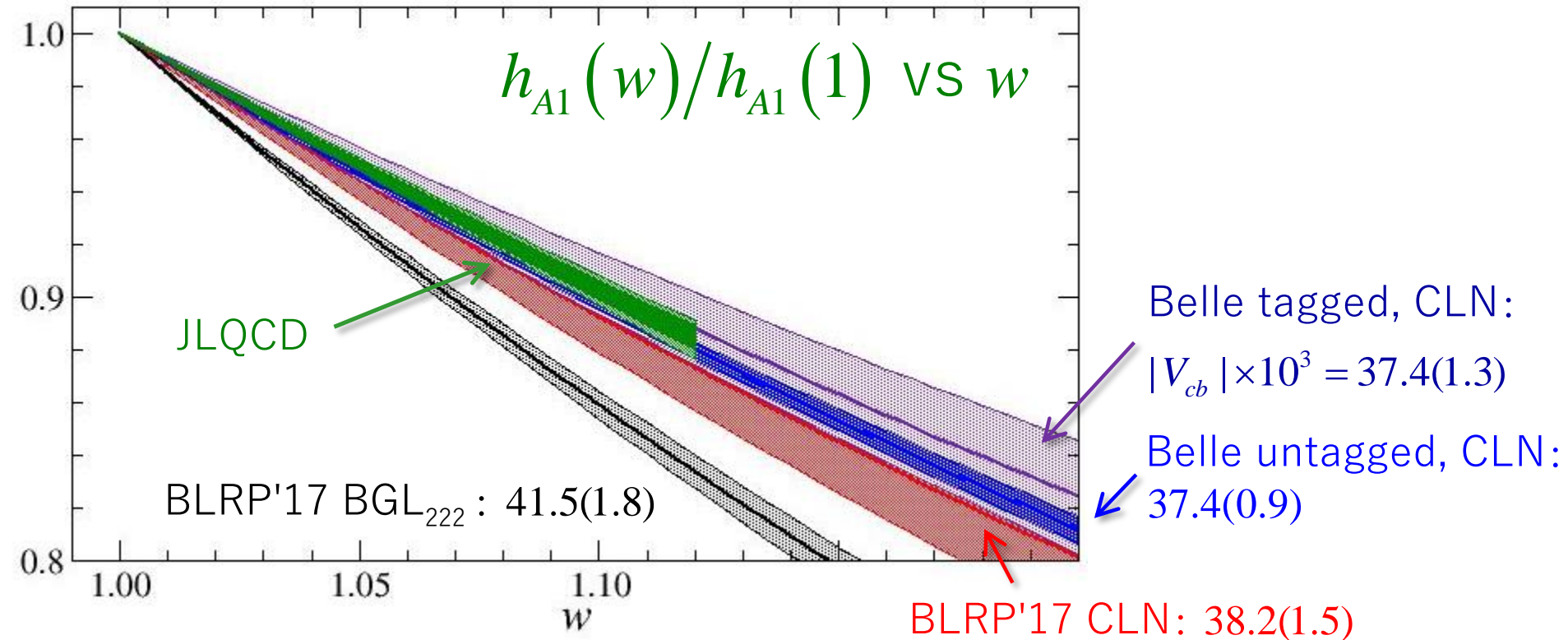
JLQCD

nicely consistent w/  
phenomenological  
analyses

reduced stat.+sys.  
error in the ratio

- good consistency among BGL / CLN / HQET / LQCD

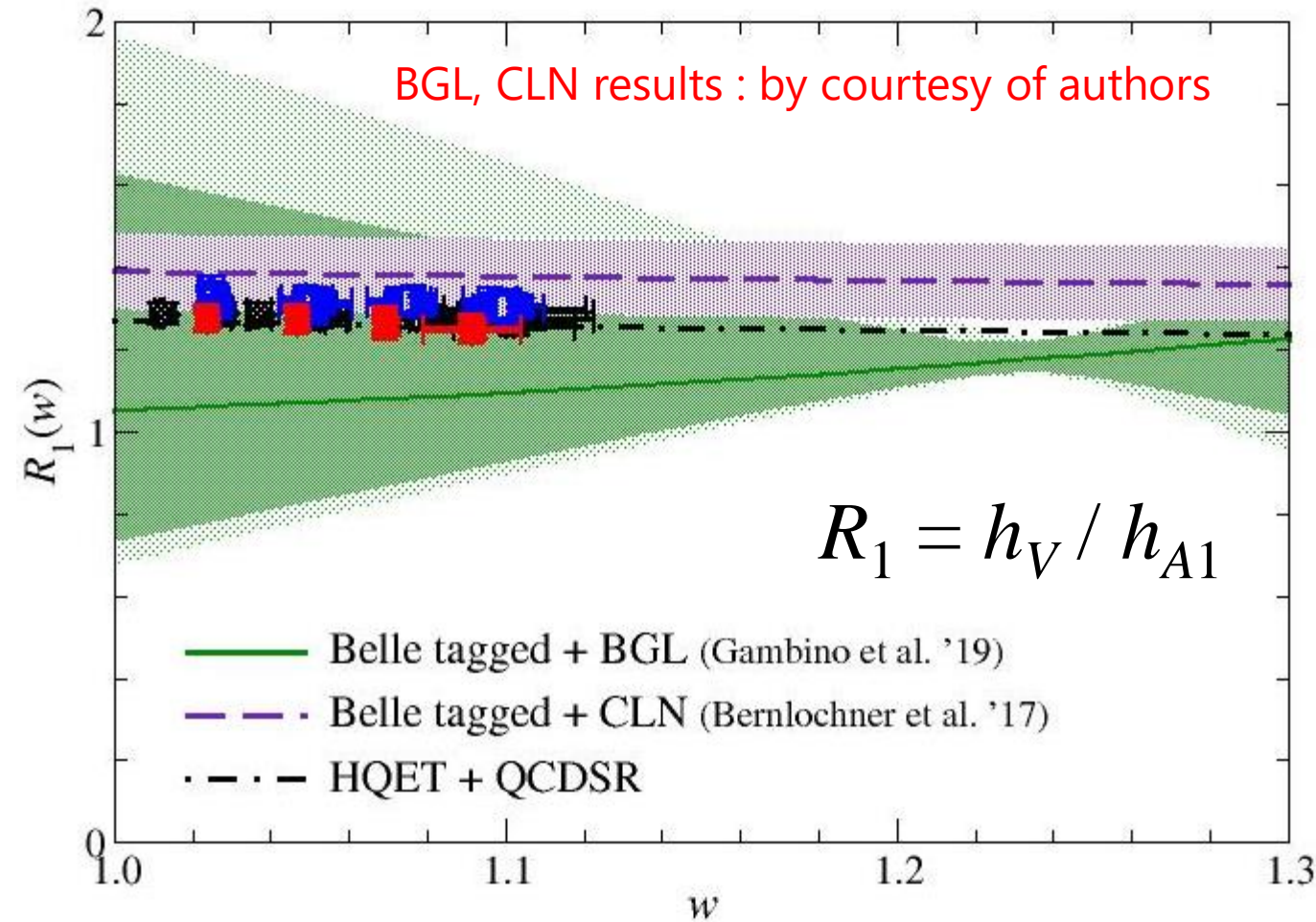
# FF shape of $h_{A1}(w)$



- '17: tension b/w CLN and BGL
- recent Belle tagged / untagged data : consistent CLN analyses
- our lattice estimate : 1<sup>st</sup> principles th. prediction, favor CLN



# FF ratio $R_1$



in 2017

BGL vs CLN, HQET

JLQCD

strongly favor  
CLN and HQET

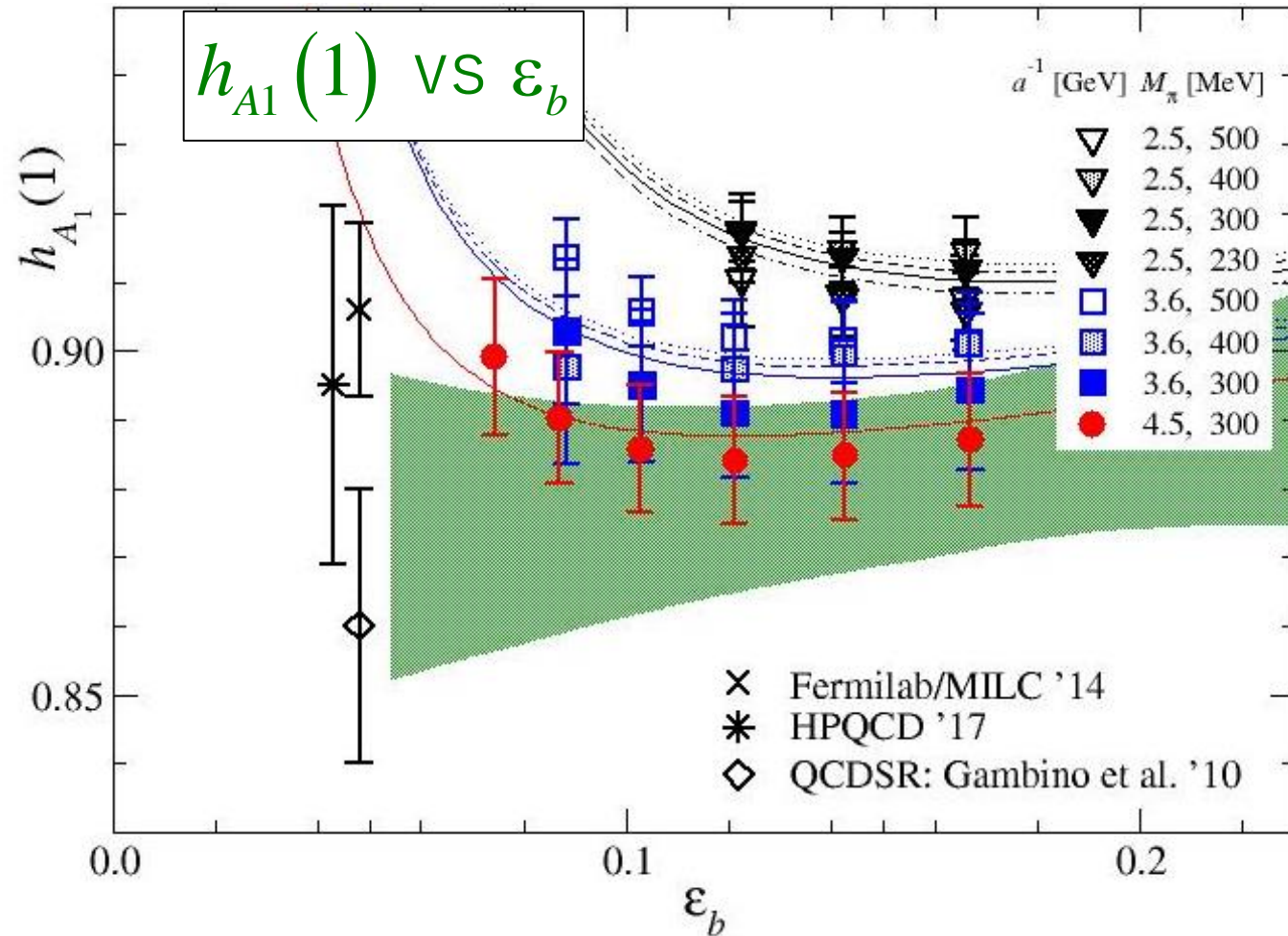
Gambino et al '19

tagged+untagged

⇒ consistent  
w/ CLN, LQCD

- systematic uncertainty of BGL analyses was underestimated
- LQCD could be helpful to stabilize "model-independent" BGL fit

# normalization $h_{A1}(1)$



HQET expansion

$$\epsilon_b = \bar{\Lambda}/M_{\eta_b}$$

$$\bar{\Lambda} = 0.5\text{GeV}$$

largest but 1-3%  
uncertainty from  
discretization

- $c_a a^2 + c_{amb} (am_b)^2$
- $c_{amb} (am_b)^2$
- $c_a a^2$

give good  $\chi^2$ 's

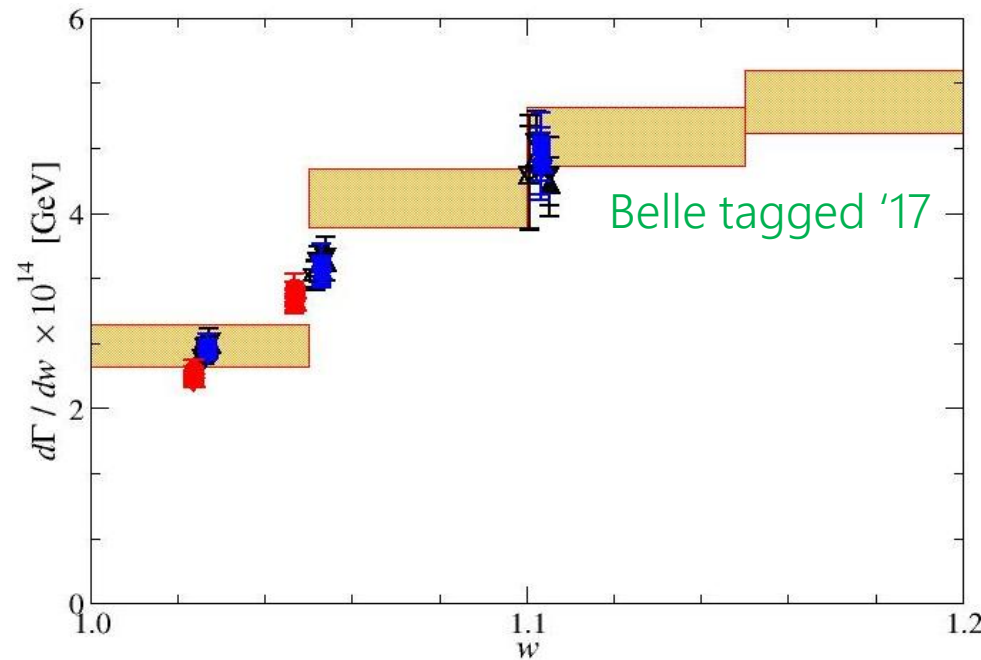
- closer to QCDSR  $\Rightarrow |V_{cb}| \times 10^3 = 38.4 - 41.1 |_{\pm 1\sigma} \Leftrightarrow 42.2(0.8)$  incl
- consistent both with previous exclusive and inclusive @  $w=1$

# towards resolution of $|V_{cb}|$ tension

collaboration w/ Belle using data in full  $w$  region

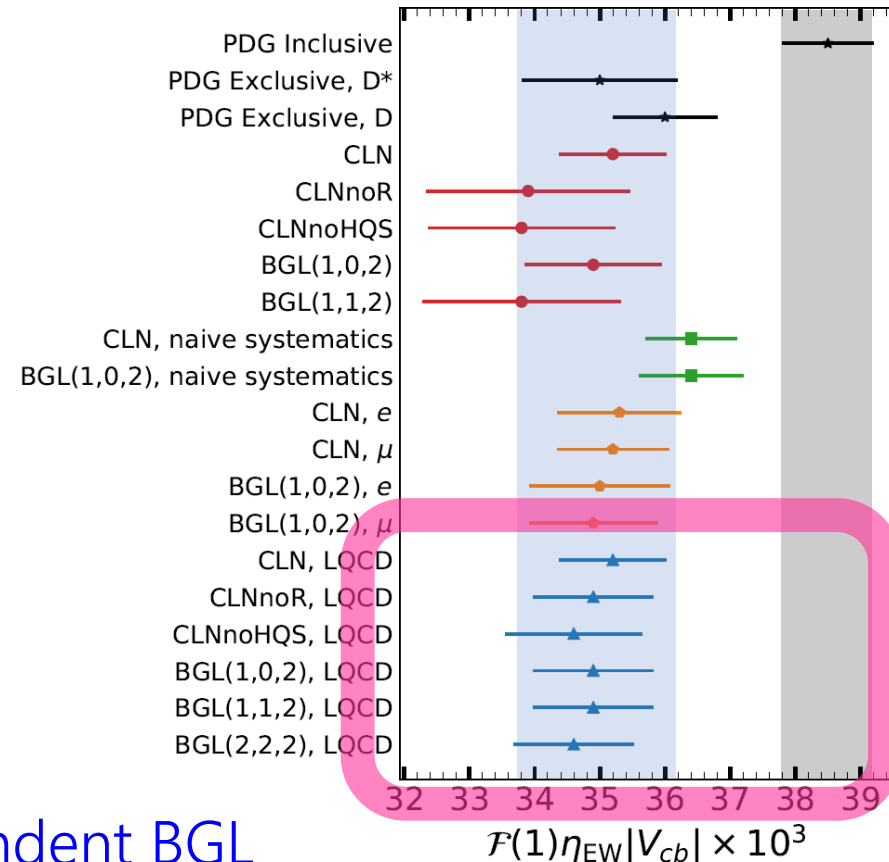
in full  $w$  region

Ferlewicz et al., '20



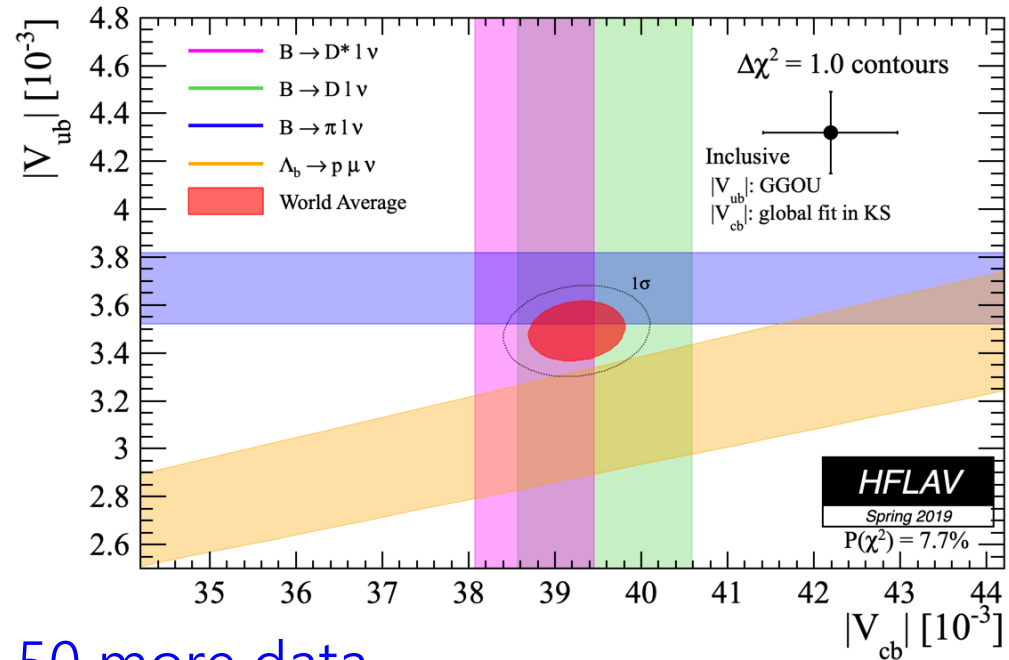
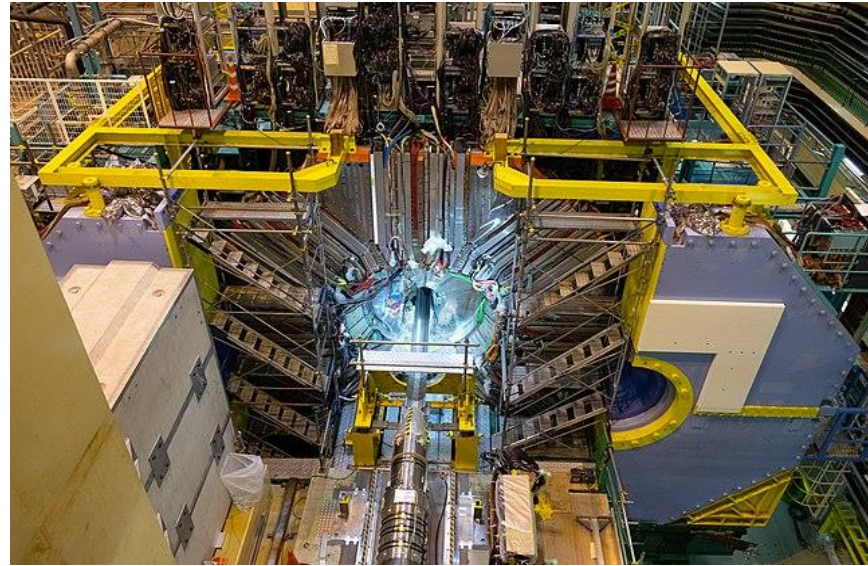
lattice data not only  $w=1$  but  $w \neq 1$

⇒ stable fit even w/ model-independent BGL



in progress

# with Fugaku and Belle II



- physics run 2019 – 2031  $\Rightarrow$  x 50 more data

- $B \rightarrow D^{(*)} \ell \nu$  : slight improvement by good control of systematics  
[ systematics cancel in the ratio  $B \rightarrow D^{(*)} \tau \nu / B \rightarrow D^{(*)} \{e, \mu\}$  ]

- $B \rightarrow \pi \ell \nu$  : big improvement in next 5 years

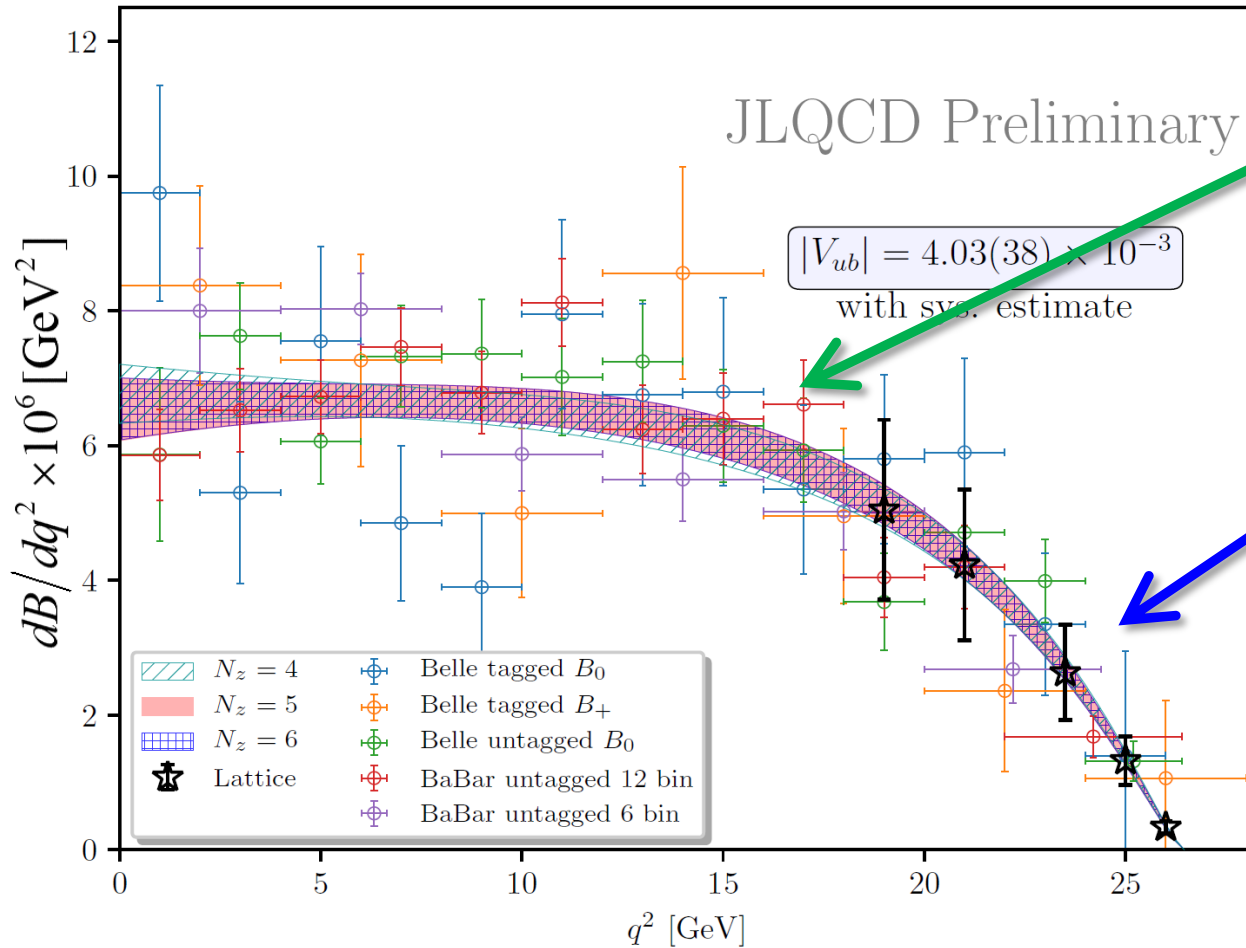
- expect progress for  $|V_{cb}|$  tension, but how about  $|V_{ub}|$ ???

**$B \rightarrow \pi \ell \nu$  to resolve  $|V_{ub}|$  tension / new physics search**



# current status

e.g. differential branching fraction



experimental data

- Belle, BaBar  $\sim 20\%$
- Belle II  $\leq 5\%$  (-2031)

our lattice QCD

- Priority Issue 9  
 $\sim 20\%$

Program for Promoting Researches : 10% in a few years

# Summary

$B \rightarrow D^* \ell \nu$  decays in relativistic lattice QCD

- promising probe of new physics, but a tension in  $|V_{cb}|$
- all form factors @ zero and non-zero recoils
  - relativistic lattice QCD w/ chiral symmetry :  $\mathcal{O}(a)$ , NPR
- indication to  $|V_{cb}|$  determination
  - CLN w/ HQET can reasonably describe FF ratios
  - LQCD would be helpful to obtain a stable BGL analysis
- on-going analysis
  - collaborative analysis with Belle
  - BSM form factors for new physics search (need NPR)